



1 Strengthening tropical influence on heat generating circulation over 2 Australia through spring

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10 Abstract

11 Extreme maximum temperatures during Australian spring can have deleterious impacts on a
12 range of sectors from health to wine grapes to planning for wildfires, but are relatively
13 understudied compared to spring rainfall. Spring maximum temperatures in Australia have
14 been rising over recent decades, and, as such, it is important to understand how Australian
15 spring maximum temperatures develop. Australia's climate is influenced by variability in the
16 tropics and extratropics, but some of this influence impacts Australia differently from winter
17 to summer, and, consequently, may have different impacts on Australia as spring evolves.
18 Using linear regression analysis, this paper explores the atmospheric dynamics and remote
19 drivers of high maximum temperatures over the individual months of spring. We find that
20 the drivers of early spring maximum temperatures in Australia are more closely related to
21 low-level wind changes, which in turn are more related to the Southern Annular Mode than
22 variability in the tropics. By late spring, Australia's maximum temperatures are
23 proportionally more related to warming through subsidence than low-level wind changes,
24 and more closely related to tropical variability. This increased relationship with the tropical
25 variability is linked with the breakdown of the subtropical jet through spring and an
26 associated change in tropically-forced Rossby wave teleconnections. However, much of the
27 maximum temperature variability cannot be explained by either tropical or extratropical
28 variability. An improved understanding of how the extratropics and tropics projects onto the
29 mechanisms that drive high maximum temperatures through spring may lead to improved
30 sub-seasonal prediction of high temperatures in the future.



31 1. Introduction

32 Anomalously high Australian spring (September-October-November) maximum
33 temperatures can be highly impactful. High temperatures may negatively impact health due
34 to a lack of acclimatisation (e.g. Nairn and Fawcett, 2014), and agriculture by changing
35 growing season length and crop yields (Cullen et al., 2009; Jarvis et al., 2019; Taylor et al.,
36 2018). Hotter and drier spring conditions have been linked to an earlier start to (Dowdy,
37 2018) and preconditioning of (Abram et al., 2021) the summer fire season. The trend toward
38 higher temperatures over recent decades (Collins et al., 2013), means that anomalous high
39 maximum temperatures may occur more often (e.g. Alexander and Arblaster, 2009). Several
40 recent springs have already exceeded historic temperature records, with some spring
41 months breaking records that were set only the previous year (Arblaster et al., 2014; Gallant
42 and Lewis, 2016; Hope et al., 2015; McKay et al., 2021). Much of this observed anomalous
43 heat has been attributed to the background global warming trend (Arblaster et al., 2014;
44 Gallant and Lewis, 2016; Hope et al., 2015; Hope et al. 2016). However, gaps remain in our
45 understanding of what drives anomalous high maximum temperatures in Australia during
46 spring, and particularly on the monthly time-scale that some of these heat events occurred
47 over. As the globe continues to warm, a better understanding of what makes a spring
48 month in Australia hot today will lead to greater resilience against extreme heat in the
49 future.

50

51 High spring temperatures have been linked with several remote modes of variability in the
52 tropics and extratropics. In the tropics, the positive phases of El Niño Southern Oscillation
53 (ENSO) in the tropical Pacific and the Indian Ocean Dipole (IOD) in the tropical Indian Ocean
54 are the strongest drivers of high maximum temperatures in Australia in spring, particularly
55 in the south and east (Power et al., 1998; Jones and Trewin, 2000; Saji et al., 2005; Min et
56 al., 2013; White et al., 2014). Many more studies focus on the ENSO and IOD relationships
57 to drier spring conditions (Nicholls et al., 1989; Meyers et al., 2007; Ummenhofer et al.,
58 2009; Risbey et al., 2009a; Watterson, 2010; Cai et al., 2011; Min et al., 2013; Pepler et al.,
59 2014; McIntosh and Hendon, 2018; Watterson, 2020) and to more extreme spring fire
60 weather (Harris and Lucas 2019; Marshall et al. 2021). While ENSO and the IOD co-vary
61 significantly in austral spring (e.g. Meyers et al., 2007), they can occur independently (e.g.



62 Risbey et al 2009a). Further, the IOD's influence on Australia's temperature peaks around
63 SON (Saji et al., 2005) compared to around NDJ (November-December-January) for ENSO
64 (Jones and Trewin, 2000). It can be useful to look at a single index that describes the large-
65 scale tropical SST variability's influence on Australia, such as the tropical tripole index (TPI)
66 (Timbal and Hendon, 2011). While other tropical modes of variability, such as the Madden-
67 Julian Oscillation (MJO), also influence Australia's spring maximum temperatures (e.g.
68 Wheeler and Hendon, 2004; Wheeler et al., 2009; Marshall et al., 2014), we focus on the
69 tropical SST-driven influence on Australia's spring climate.

70

71 Variability in the extratropics is also linked to high temperatures in Australia. The negative
72 phase of the Southern Annular Mode (SAM), the primary mode of variability in the
73 extratropics, (Hendon 2007; Risbey et al., 2009a; Min et al., 2013; Marshall et al., 2012,
74 Hendon et al., 2014; Fogt and Marshall, 2020) drives hotter and drier Australian spring
75 conditions, and to more extreme spring fire weather (Marshall et al. 2021). SAM generally
76 varies at a higher frequency than ENSO or the IOD, however, SAM also has lower frequency
77 variations. On a seasonal timescale, El Niño promotes negative SAM, particularly during the
78 warmer months (L'Heureux and Thompson, 2006; Hendon et al., 2007; Lim et al., 2016; Lim
79 et al., 2019a). Polar stratospheric weakening during austral spring (sometimes associated
80 with sudden stratospheric warming) can also sustain negative SAM (Lim et al., 2018)) and
81 higher Australian maximum temperatures from late spring (Lim et al., 2019b). As with ENSO
82 and the IOD, more studies focus on how SAM influences Australian rainfall than
83 temperature, particularly when examining the teleconnection pathway. While low rainfall
84 correlates well with high maximum temperatures (Simmonds, 1998; Jones and Trewin,
85 2000; Timbal et al., 2002; Hope and Watterson, 2018), there is a gap in our understanding of
86 how both tropical and extratropical modes of variability impact spring maximum
87 temperature.

88

89 Anomalous high geopotential height (or, synonymously, anticyclonic vorticity) over
90 southern Australia is associated with spring high maximum temperatures in Australia (Hope
91 et al., 2015; Gallant and Lewis 2016; McKay et al., 2021). While ENSO, the IOD, and the
92 tropical TPI also promote anomalously high geopotential height, it forms further to the
93 south of Australia (e.g. Cai et al., 2011; Timbal and Hendon, 2011; McIntosh and Hendon,



94 2018). SAM's negative phase is characterised by an equatorward shift of the eddy-driven jet
95 and bands of anomalously low and high geopotential height in the mid- and high-latitudes
96 respectively (Thompson and Wallace, 2000; Fogt and Marshall, 2020). The altered
97 atmospheric flow associated with the drivers can reduce rainfall, including by defecting
98 cooling rain-bearing systems (e.g. Jones and Trewin 2000; Hendon et al., 2007; Pepler et al.,
99 2014; van Rensch et al., 2019) away from Australia (Cai et al., 2011; Risbey et al., 2009b;
100 McIntosh and Hendon, 2018; Hauser et al., 2020). Anomalous heat and dry is also
101 associated with other mechanisms such as increased subsidence and insolation (Hendon et
102 al., 2014; Lim et al., 2019b; Pfahl et al., 2015; Quinting and Reeder, 2017; Suarez-Gutierrez
103 et al., 2020) or heat advection (Jones and Trewin, 2000; Boschat et al., 2015; Gibson et al.,
104 2017). Understanding the differences between the extratropical and tropical forcing behind
105 some of these heat mechanisms is a goal of this paper.

106

107 The mechanisms and atmospheric circulation patterns associated with heat and connections
108 to remote drivers may also vary through spring. McKay et al. (2021) noted that the
109 relationship with the southern Australian upper-anticyclone and maximum temperature is
110 weaker in September than November, and suggested that the anticyclone had greater
111 influence from the tropics in later spring. The impact of SAM in the extratropics on
112 Australia's temperature reverses from winter to spring (Hendon et al., 2007; Risbey et al.,
113 2009a; Marshall et al., 2012; Min et al., 2013; Hendon et al., 2014; Fogt et al., 2020) as the
114 mean zonal winds change with the seasons (Hendon et al. 2007) and the Indo-Pacific
115 subtropical jet (STJ) weakens (Bals-Elsholz et al., 2001; Koch et al., 2006; Ceppi and
116 Hartmann, 2013, Gillett et al., 2021) so that a negative SAM phase enhances subsidence
117 over subtropical Australia into the warmer months (Hendon et al., 2014). The IOD and ENSO
118 teleconnection pathways over the Indian Ocean toward Australia also change from winter
119 to spring (Cai et al., 2011). This change may relate to the strength of the winter STJ, as it
120 should prevent direct propagation of Rossby waves between the tropics and extratropics
121 (e.g. Hoskins and Ambrizzi, 1993). McIntosh and Hendon (2018) proposed that transient
122 eddy-feedbacks generate a secondary wave source south of the winter STJ in response to
123 IOD forcing. In spring, the STJ weakens sufficiently to allow for direct Rossby wave
124 propagation from the tropical Indian Ocean. However, McKay et al. (2021) suggested that



125 the STJ may not weaken sufficiently in September to allow direct Rossby wave propagation,
126 and that teleconnection pathways may be different on a monthly timescale as a result.

127

128 Teleconnections driven by large-scale remote modes of variability can precondition
129 Australia toward hotter and spring conditions (e.g. Hurrell et al., 2009), but cannot
130 guarantee a hot month or season will eventuate. Even the strongest El Niño events may not
131 result in the canonical dry and warm conditions expected (van Rensch et al., 2019; Hauser et
132 al., 2020). Further, differences in how those modes of variability influence Australia
133 between winter-spring-summer and the differences between spring-average atmospheric
134 circulation highlight that there is more to understand in how maximum temperatures evolve
135 through spring months. Filling the gap between weather and seasonal time-scales is an
136 ongoing area of research that can lead to improved sub-seasonal forecasting (Meehl et al.,
137 2021). Given the increasing likelihood of future extreme heat events occurring through
138 spring, it is imperative to understand any differences that may exist in how heat develops,
139 and links to varying influences from the extratropics to tropics. The reanalysis datasets,
140 Rossby wave and statistical analysis methods are described in Section 2. An overview of how
141 Australian spring maximum temperatures are related to circulation and large-scale
142 variability is in Section 3. In Section 4 the variation of these relationships through the
143 months of spring are assessed and Section 5 describes how the drivers influence the
144 mechanisms that promote high monthly maximum temperature. Discussion and conclusions
145 are provided in Section 6

146

147 [2. Methods and data](#)

148 [2.1 Indices and datasets](#)

149 All circulation variables for September, October, November monthly-averaged data are
150 taken from the ECMWF's Reanalysis 5 (ERA5) (Hersbach et al., 2020) available from the
151 Copernicus Climate Change Service (C3S, 2017) on a 0.25° grid from 1979 to 2019. Low-level
152 circulation is diagnosed using 850hPa horizontal wind and mean sea level pressure (MSLP).
153 Mid-tropospheric vertical motion is represented by 500hPa omega. Upper-level circulation
154 is represented by 200hPa geopotential height (200Z). 200hPa horizontal winds are used for



155 Rossby wave analysis. Similar results were found using ERA-Interim reanalysis (Dee et al,
156 2011) and the JRA-55 from the Japan Meteorological Agency (2013) (not shown).
157
158 Australian monthly-averaged daily maximum temperature data for 1979 to 2019 is taken
159 from the Australian Water Availability Project (AWAP) (Jones et al., 2009) analyses, available
160 on a 0.05° resolution grid.
161
162 Monthly sea surface temperature (SST) is taken from NOAA Extended Reconstructed Sea
163 Surface Temperature (ERSST V5; Huang et al., 2017)
164
165 The impacts of SAM on Australia's climate shows some sensitivity to the method used to
166 calculate the SAM index (e.g. Risbey et al., 2009a). To ensure consistency between the other
167 indices and circulation variables, we calculate SAM as the difference between the
168 standardized zonal means of ERA5 MSLP anomalies at 60°S and 40°S (Gong and Wang,
169 1999).
170
171 The tropical TPI (Timbal and Hendon, 2011) is defined as the difference in SST averaged over
172 a parallelogram located over the Maritime Continent (0°-20S, 90°-140E at the equator
173 shifted to 110°-160°E at 20°S) from SST averaged and summed over two regions in the
174 tropical Indian Ocean (10°N to 20°S, 55° to 90°E) and tropical Pacific Ocean (a trapezium
175 that extends from 15°N to 15°S, 150°E to 140°W in the north and 180°E to 140°W in the
176 south). ENSO is described using the Niño3.4 index (averaged SST anomalies over 5°N-5°S,
177 170°E-120°W) and the IOD using the dipole mode index (DMI; the difference between the
178 SST anomalies averaged over 10°S-10°N, 50°-70°E and 10°S-0°, 90°-110°E; Saji et al., 1999).
179
180 To highlight the influence of interannual variability, the 1981-2010 climatological mean is
181 removed from each month, and the data is linearly detrended before analysis.
182



183 2.3 Rossby wave analysis

184 We use wave activity flux (WAF) at 200hPa to trace Rossby wave group propagation and to
185 identify source and decay regions that influence the atmospheric circulation patterns.

186 Following Takaya and Nakamura (2001), we calculate WAF as:

187

$$\begin{aligned} 188 \quad WAF = p \cos \phi & \left\{ \frac{U}{a^2 \cos^2 \phi} \left[\left(\frac{\partial \psi'}{\partial \lambda} \right)^2 - \psi' \frac{\partial^2 \psi'}{\partial \lambda^2} \right] \right. \\ 189 & + \frac{V}{a^2 \cos \phi} \left[\frac{\partial \psi'}{\partial \lambda} \frac{\partial \psi'}{\partial \phi} - \psi' \frac{\partial^2 \psi'}{\partial \lambda \partial \phi} \right] \frac{U}{a^2 \cos \phi} \left[\frac{\partial \psi'}{\partial \lambda} \frac{\partial \psi'}{\partial \phi} - \psi' \frac{\partial^2 \psi'}{\partial \lambda \partial \phi} \right] \\ 190 & \left. + \frac{V}{a^2} \left[\left(\frac{\partial \psi'}{\partial \phi} \right)^2 - \psi' \frac{\partial^2 \psi'}{\partial \phi^2} \right] \right\} \end{aligned}$$

191

192 where p is the pressure (200hPa) scaled against 1000hPa, U and V are the climatological
193 zonal and meridional wind speed magnitudes, a is the radius of the earth, (ϕ, λ) are latitude
194 and longitude, $\psi' = Z' / f$ is the quasi-geostrophic perturbation streamfunction, Z' is the
195 200hPa geopotential height anomaly obtained through regression onto maximum
196 temperature or climate driver indices, $f = 2\Omega \sin \phi$ is the Coriolis parameter with the
197 Earth's rotation Ω . WAF is not plotted within 10° of the equator.

198

199 WAF propagates in the direction of quasi-stationary Rossby wave group velocity, and
200 regions of divergence or convergence of WAF correspond to zones of Rossby wave sources
201 or sinks respectively.

202

203 Total stationary Rossby wave wavenumber (e.g., Hoskins and Karoly 1981) is defined as:

$$204 \quad K_S = \sqrt{\frac{\beta - U_{yy}}{U}}$$

205 where $\beta - U_{yy}$ is the meridional gradient of mean-state absolute vorticity at 200hPa. WAF
206 should refract toward regions of higher K_S and either reflect or evanesce on regions of $K_S < 0$,
207 such as in the STJ where the curvature of the flow (U_{yy}) can become larger than the
208 planetary vorticity gradient (β) (e.g Barnes and Hartmann, 2012; Li et al., 2015 a,b)

209



210 2.4 Statistical analysis

211 Linear, partial, and multi-linear regression and Spearman's ranked correlation are used to
212 assess the relationships between Australian maximum temperature, atmospheric circulation
213 and the tropics and extratropics. Due to the large decorrelation length scales, Australian-
214 average maximum temperature variability is representative of all but far north Australia's
215 spring and spring-monthly maximum temperatures (Sup. Fig. 1). Statistical significance is
216 calculated at the 95% confidence level using Student's (1908) t-test using 39 (41 years - 2)
217 degrees of freedom. Pattern correlation is used to compare regression patterns.

218

219 3. Spring-season maximum temperatures - circulation patterns and 220 associations with drivers

221

222 We start by giving an overview of the spring-seasonal relationships between average
223 Australian austral spring maximum temperature and lower- and upper-level atmospheric
224 circulation (Fig 1a,b). Barotropic cyclones appear to the southwest and southeast of
225 Australia, occurring in both the lower- and upper-level circulation regressions (Fig. 1a-b) and
226 noted during recent extreme spring heat events (Gallant and Lewis, 2016; Hope et al., 2016;
227 McKay et al., 2021). Weak anticyclonic low-level winds are found over Australia, as well as
228 sinking motion across the eastern half of the continent. An upper-level anticyclone sits over
229 southern Australia, with the wave activity flux predominantly propagating from the
230 subtropical Indian Ocean, through the anticyclone and into the subtropical Pacific Ocean.

231

232 We now compare the atmospheric patterns associated with spring maximum temperature to
233 those associated with large scale modes of variability. The spring-average atmospheric
234 circulation patterns associated with the remote drivers of variability are calculated via linear
235 regression onto each standardised index. Note that the TPI and SAM indices have been
236 multiplied by negative one to present positive associations with high temperatures. The
237 pattern for SAM (x-1) shows elongated barotropic low and high anomalies lie in the middle
238 and high latitudes respectively (Fig. 1c-d), with upper-level cyclonic nodes to the southeast



239 and southwest of Australia. Negative SAM is associated with high maximum temperatures
 240 through much of subtropical, and particularly eastern, Australia (Fig. 1e).
 241

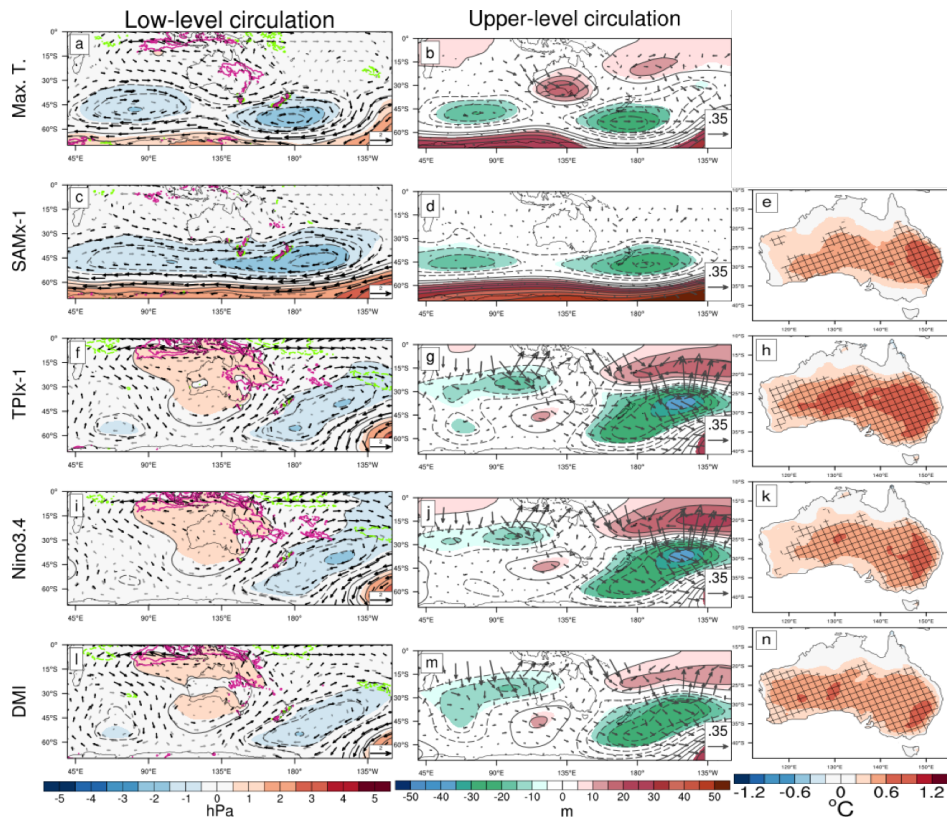


Figure 1. Linear regressions of spring standardised weighted area-averaged Australian maximum temperature (a-b), SAMx-1 (c-e), tropical TPIx-1 (f-h), Niño3.4 (i-j) and DMI (j-n) onto low-level circulation (left column), upper-level circulation (middle column) and Australian maximum temperatures (right column). Low-level circulation is represented by anomalous mean sea level pressure (hPa) (black and filled contours), 850hPa wind vectors (ms⁻¹) and 500hPa omega (hPas⁻¹) contours from -0.02 to 0.02hPas⁻¹ in steps of 0.01 hPas⁻¹ (magenta contours are positive; downward motion) and cyan contours are negative; upward motion, and the zero contour is not plotted). Upper-level circulation is represented by 200hPa geopotential height (black and filled contours) and wave activity flux vectors (m²s⁻²). Filled contours, bold wind vectors, cross-hatching, and all vertical motion contours are significant at the 95% confidence level using a Student's t-test with 39 independent

242
 243 The tropical modes, represented by Niño3.4 (Fig 1i-k), the DMI (Fig. 1 l-n), and tropical TPI
 244 (x-1) (Fig 1f-h)) are also associated with spring Australian maximum temperature anomalies.
 245 Each mode generates an apparent Rossby wave pattern that arcs from the tropical Indian



246 Ocean to promote anomalous high geopotential height south of Australia, consistent with
247 earlier studies (e.g. Cai et al., 2011; Timbal and Hendon, 2011; McIntosh and Hendon, 2018).
248 Each regression also shares anomalous high surface pressure over Australia, sinking motion
249 in the east, cyclonic nodes to the southwest and east of Australia, and elongated upper-level
250 cyclones in the subtropical Indian Ocean. These similarities are likely the result of the strong
251 co-variability between the IOD and ENSO (e.g. Meyers et al., 2007; Risbey et al., 2009a).
252 However, the IOD has a stronger low-level cyclone to the southeast and a poleward
253 extension of the subtropical Indian Ocean cyclone that sets a subtly different wave train
254 from around 50°S, 60°E that is poleward that generated by ENSO. The positive IOD is also
255 associated with high maximum temperatures across a broader region of southern and
256 western Australian than is El Niño. The tropical TPI ($x-1$) is a blend of both Niño3.4 and DMI
257 circulation patterns and has a strong relationship with Australian spring maximum
258 temperatures across all but northern Australia.

259

260 Given the similarities and connections between ENSO and IOD teleconnections, we use the
261 tropical TPI to represent the large-scale influence of the tropics. SAM is used to represent
262 the influence of the extratropics. Statistical models of Australian weighted area-averaged
263 spring maximum temperatures reconstructed through multilinear regression using either
264 Niño3.4, DMI and, SAM or the tropical TPI and SAM as the predictors explains 32% and 34%
265 of maximum temperature variability respectively (sup. Fig. 2).

266

267 We next compare the atmospheric circulation associated with monthly high maximum
268 temperatures to that with the large-scale modes of variability through the individual months
269 of spring. To ensure that we are assessing the influence of the tropics and extratropics
270 separately, we use multi-linear regression onto the monthly circulation variables.

271

272 4. Monthly circulation patterns and associations with drivers

273 The regression of monthly Australian maximum temperature onto the lower- and upper-
274 level atmospheric circulation is displayed in Figures 2a-c and 3a-c respectively for
275 September, October and November. The multi-linear regression onto the standardised
276 monthly indices of SAM ($x-1$) (Figs. 2d-f and 3d-f) and tropical TPI ($x-1$) (Figs. 2h-j and 3h-j).



277 At first glance, these monthly circulation patterns are broadly similar to the spring-average
278 regression patterns. However, the details of the circulation patterns change as the months
279 progress, suggesting that different processes are important for heat development through
280 spring.
281

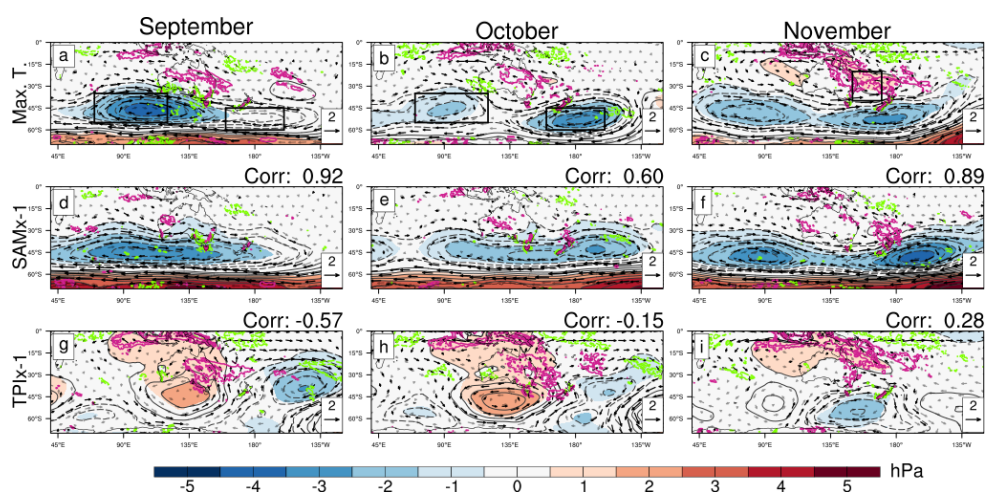


Figure 2. Regressions onto low-level circulation, as in Fig. 1, except for September (left column), October (middle column) and November (right column). Standardised areal-averaged Australian maximum temperature is linearly regressed onto low-level circulation (ta-c) and SAMx-1 (d-f) and tropical TPIx-1 (g-i) are multi-linear regressions onto low-level circulation.

Pattern correlation between the maximum temperature MSLP regressions and the SAM and tropical TPI regressions calculated over 5°S-70°S; 70°E-170°E are written in the top right of each SAM or tropical TPI regression.

The boxes (a-c) show key low-level circulation features identified as being important for maximum temperature development: The southwest cyclone (SWC) 35°S-55°S; 70°-120°E; southeast cyclone (SEC) 45°S-60°S; 160°-200°E and Tasman Sea high (TSH) 20°S-40°S; 150°-170°E

282
283 The most obvious change in atmospheric circulation through the months is in the low-level
284 flow across Australia, particularly generated by the barotropic cyclones southwest (SWC) or
285 southeast (SEC) of Australia (boxes in Fig 2a-b). Weak low-level anticyclonic flow around the
286 Tasman Sea (box in Fig. 2c) also contributes to the northerly flow over eastern Australia in
287 November in particular (Fig. 2c). Tasman Sea anticyclonic blocking patterns have previously



288 been linked to anomalously warm conditions (Marshall et al., 2014), but here appear to only
289 contribute to high maximum temperatures in November. The SWC and SEC vary in
290 geographic shape and strength through the months. The SWC dominates in September but
291 weakens through October and November, whereas, the SEC is missing in September but is
292 strong in October and November. Similar cyclones appear in the monthly SAM ($x-1$)
293 regressions (Figs. 2 d-f, 3d-f) and the Australian-region MSLP correlates strongly with that
294 associated with high Australian temperature (top-right of Fig 2d-f). Rather than cyclones in
295 September and October the TPI ($x-1$) is associated with a barotropic anticyclone south of
296 Australia that directs southerly low-level wind across eastern Australia (Figs. 2 h-j); a pattern
297 that would be associated with cooler conditions. The September and October TPI ($x-1$) MSLP
298 pattern actually anti-correlate with that associated with high Australian maximum
299 temperatures (top-right Fig. 2 g-i). It is not until November that we see a barotropic cyclone
300 to the southeast of Australia associated with the TPI ($x-1$). So, for the majority of spring
301 negative TPI-forced low-level atmospheric circulation appears to counter high maximum
302 temperatures, despite the overall positive relationship in spring (Fig. 1h).

303

304 The anomalous southern Australian upper-anticyclone (SAA) from the spring pattern
305 appears is also associated with high maximum temperature in each of the individual spring
306 months (Fig. 3a-c), but its location shifts eastward across Australia through spring. The
307 boxed region was chosen to match earlier studies (Gallant and Lewis, 2016; McKay et al.,
308 2021), but best matches the November position, likely contributing to the stronger
309 relationship between heat and SAA in this month (McKay et al., 2021; see also section 6).
310 The anticyclone in later spring appears to form part of a wave train from a cyclone to the
311 northwest of Australia toward the southeast cyclone. While the monthly TPI regressions
312 have anticyclones in September and October (Fig. 3g-h), they are located too far south
313 relative to Australia, as in the spring-average regression. The regressions onto SAM ($x-1$) (Fig
314 3d-f) have weak anticyclones over western Australia that are not statistically significant. It is
315 not until November that both SAM and TPI ($x-1$) (Figs 3 f, i) have an anticyclone over central-
316 east southern Australia. Both the upper- level SAM and TPI ($x-1$) regressions correlate
317 moderately with the maximum temperature regression in November, and the SAM and TPI
318 ($x-1$) anticyclones may form part of the same wave train associated with maximum
319 temperatures. However, the SAM and TPI ($x-1$) anticyclones are weaker and too far east



320 relative that associated with high maximum temperatures, such that they may not
321 contribute strongly to the SAA formation. We explore this idea further in section 6.
322

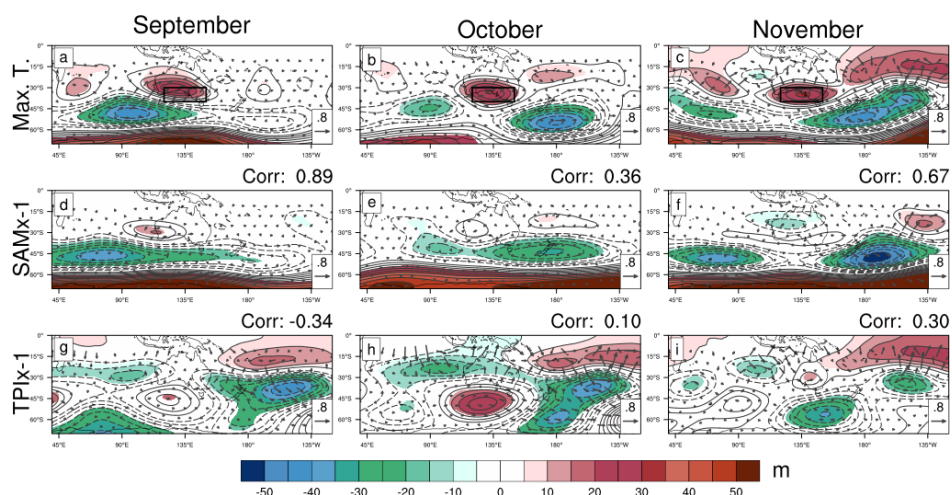


Figure 3. As with figure 2, but for upper-level circulation. The Australian-region pattern correlation between the maximum temperature Z200 regressions and SAMx-1 and TPIx-1 are in the top right of each figure. Boxed area (a-c) highlights the southern Australian anticyclone (SAA; 30°-40°S, 120°-150°E) that is linked with high maximum temperatures.

323
324 While the southern Australian anticyclone is not well explained by SAM or TPI (x-1) through
325 spring, much of the statistically significant 500hPa vertical motion associated high maximum
326 temperatures (green and magenta contours, Fig. 2a-c) matches that associated with TPI (Fig.
327 2h-j) and to a lesser extent SAM (Fig. 2d-f). In September, sinking motion over subtropical
328 Australia and rising motion over the southern coasts is associated with high maximum
329 temperatures. By November, the rising motion has largely vanished and the sinking motion
330 has shifted to be over eastern Australia. It was expected that the SAA would generate some
331 of the sinking motion associated with high maximum temperature, however, this vertical
332 motion does not correlate strongly with any of the key circulation features examined here
333 (Sup. Table1).
334



335 Changes in propagation of wave activity flux help explain some of the changes in the broad
336 scale circulation changes through spring. In September, WAF predominantly propagates
337 from the southwest cyclone toward the southern Australian anticyclone. In October, a
338 component of WAF also propagates from the eastern tropical Indian Ocean region. By
339 November, the tropical-component dominates the WAF and forms part of a very different
340 pattern to the previous two months; continuous WAF propagates from the far southwest
341 Indian Ocean, joins WAF propagating out of the tropical Indian Ocean and then continues
342 across the southern Australian anticyclone. The latter part of this wave train is similar to the
343 IOD teleconnection highlighted by Cai et al. (2011) in spring. The WAF associated with SAM
344 and TPI ($x-1$) also propagates from the extratropics toward the respective anticyclones in
345 September and October. While a broad region of low height in the subtropical Indian Ocean
346 is associated with the TPI, it does not appear to generate WAF that propagates into the
347 extratropics. It is not until November that WAF associated with the TPI ($x-1$), and weakly
348 with SAM ($x-1$), appears to propagate directly from the cyclone in the eastern subtropical
349 Indian Ocean through the anticyclone over southeastern Australia.

350

351 Overall, these results suggest that the circulation associated with maximum temperature
352 shifts from extratropical to tropical forcing as spring progresses. This is supported by how
353 well SAM appears to project onto the atmospheric circulation associated with maximum
354 temperatures in September, and how the TPI projects more strongly later in spring. The
355 change in WAF associated with this change suggests that there may be a blocking
356 mechanism between the tropics and extratropics generating this change.

357

358 We find qualitatively similar results if we perform the linear regressions using maximum
359 temperature averaged over sub-regions of Australia, for example southwest or southeast
360 Australia (Supplemental Fig S2).

361

362 5. Connection between subtropical jet and atmospheric circulation

363 We next explore how the subtropical jet may be influencing the WAF through the spring
364 months.

365



366 The subtropical jet (STJ) peaks in strength in winter and weakens through spring to have
367 broken down by summer (e.g. see figure 9, Ceppi and Hartmann, 2013). This gradual
368 breakdown of the STJ coincides with a decrease in the area with total stationary
369 wavenumber less than zero over southern Australia (Fig. 4), and may provide an explanation
370 for the growing relationship with the tropics and Australian maximum temperature by
371 November.

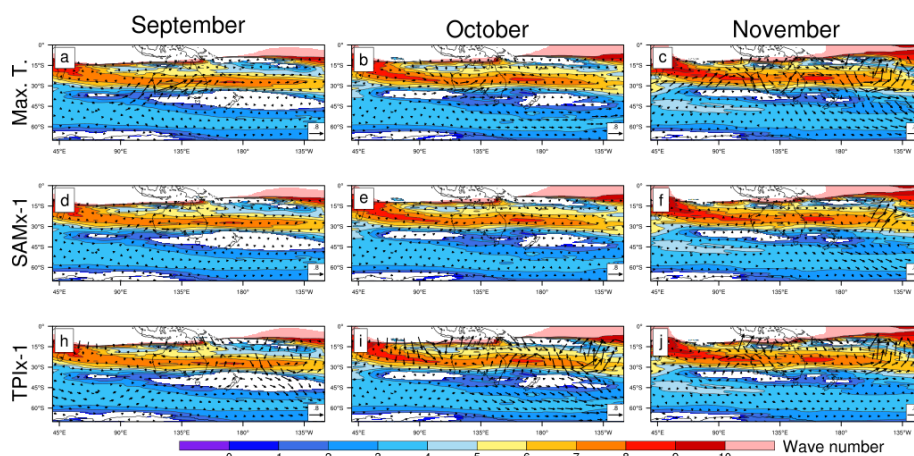


Figure 4. Total wave number (K) calculated for September, October and November. Vectors are the wave activity flux repeated from figure 3.

372
373 The wave activity flux vectors from the maximum temperature, TPI and SAM (x-1)
374 regressions in figure 3 are overlaid in figure 4 on the monthly climatological Ks associated
375 with the zonal winds. In September, the WAF associated with high maximum temperature
376 (Fig. 4a) diverges from the region of the southwest cyclone to propagate through a region of
377 low total stationary Rossby wave wavenumber, Ks, over southwest Australia and along the
378 STJ waveguide (i.e. from high to low latitudes). As the jet weakens in October (Fig. 4b) a
379 portion of WAF also diverges from the tropical Indian Ocean to dissipate on the jet's
380 equatorward flank, but mostly propagates from west to east along the STJ waveguide. Even
381 more distinctive, by November (Fig 4c), WAF propagates along the jet waveguide to a region
382 near Africa, with contributions from the tropical Indian Ocean, but does not appear to
383 propagate out of the SWC.
384
385 The increase in WAF associated with the tropical TPI (Fig 4h-j) propagating out of the
386 tropical Indian Ocean through spring appears to coincide with the STJ decay. In September



387 and October weak WAF diverging from the central southern Indian Ocean follows the eddy-
388 driven jet waveguide (region of locally higher wave number around 50°S), suggesting the
389 secondary wave source proposed by McIntosh and Hendon (2018) is important in early
390 spring. The tendency for TPI-associated WAF to form and follow this trajectory may explain
391 why the barotropic anticyclone associated with the TPI is further poleward than in the
392 regression onto Australian maximum temperature. By October more WAF is propagating
393 out of the tropical Indian Ocean along the region of high Ks and by November WAF is
394 propagating out of the extratropical Indian Ocean along the high Ks region, similar to the
395 maximum temperature-WAF. WAF generated by SAM (Fig. 4d-f) also converges toward the
396 STJ waveguide in each month.

397

398 Limits around linear Rossby wave theory (e.g. Liu & Alexander, 2007) may explain why some
399 wave activity flux cross the region of imaginary wavenumber associated with the STJ.
400 However, the majority of WAF associated with Australian maximum temperature, or with
401 the tropics or extratropics does divert to propagate along the jet, as expected. While the
402 breakdown of the STJ through spring may help explain the change in teleconnection
403 pathways of the TPI toward Australia, the STJ consistently acts as a waveguide toward
404 Australia.

405

406 We now look more closely into how the drivers, circulation features, and heat mechanisms
407 relate to each other and how that results in higher Australian maximum temperatures.

408

409 [6. Mechanisms and drivers of monthly maximum temperatures through spring](#)

410 As with the atmospheric circulation regressions, the relationships between Australian
411 maximum temperature and SAM and TPI ($x-1$) evolve through the spring months. In
412 September, negative SAM (Fig. 5a) is associated with a broad area of high maximum
413 temperature over subtropical Australia, that contracts in October and November (Figs. 5 b-
414 c). Conversely, the relationship with negative TPI and maximum temperature is weaker early
415 in spring, with statistically significant high temperatures confined to the west and east, and
416 cool temperatures in the far north in September (Fig. 5d). The TPI's relationship with high
417 maximum temperature broadens and strengthens in October and covers the majority of



418 Australia by November (Figs. 5 e-f). Overall, these monthly relationships give the impression
419 of a transition from extratropical to tropical drivers becoming more influential over
420 Australian temperatures that is broadly consistent with the apparent change in atmospheric
421 circulation through spring.

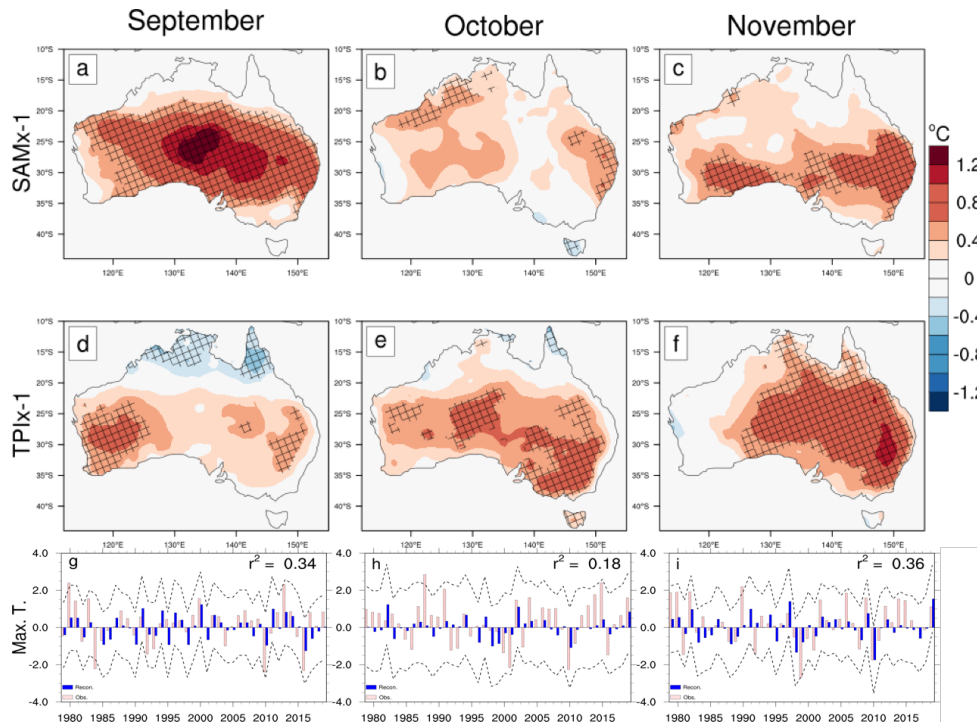


Figure 5. Multilinear regression coefficients (°C) of Australian maximum temperature regressed onto standardised timeseries of the SAMx-1 (a – c) and the tropical TPI x-1 (d – f) for September, October and November over the years 1979 to 2019.

Reconstructions (blue bars) of September, October and November (i-k) Australian area-averaged maximum temperature from standardised time series of SAM and tropical TPI indices. Observed values are in blue. The dashed line shows the 95% prediction interval computed as ± 1.96 standard error and the variance explained (r^2) of the model is in the top right of each figure.

422
423 Using the standardised SAM and TPI time series as predictors in a regression model to
424 reconstruct the monthly Australian-averaged maximum temperature anomalies (Figs 4g-i)
425 explains only between 18 and 36% Australian maximum temperature variance (r^2) through
426 spring. The model does not substantially improve if it is calculated over southeast or
427 southwest Australia, or if using Niño3.4 or DMI as predictors instead of the tropical TPI (Sup.
428 Fig.4).



429

430 To explore how the atmospheric circulation relates to some of the mechanisms that develop
431 heat through spring, we first compose indices of the key circulation features discussed in
432 section 4. Weighted area-averages of mean-sea level pressure (multiplied by negative one)
433 over the southwest and southeast cyclones (SWC and SEC) and 200hPa geopotential height
434 over the southern Australian anticyclone (SAA) for each spring month. See Figs. 2a,b and 3a-
435 c for regions. Creating a statistical model of Australian-averaged monthly spring maximum
436 temperatures from these circulation features (Fig. 6a-c) explains consistently higher
437 maximum temperature variance (around 60%) than did the model from the indices of
438 tropical and extratropical large-scale modes of variability. Further, despite the changes in
439 the features' geographic shape, strength and position across the spring months in Fig. 2, the
440 majority of maximum temperature across Australia is well explained by at least one of these
441 features at all times through spring (Sup. Fig. 6). We next explore how these MSLP or
442 200hPa geopotential height features relate to the low-level westerly or northerly winds and
443 vertical motion and how that relates to high maximum temperature development.

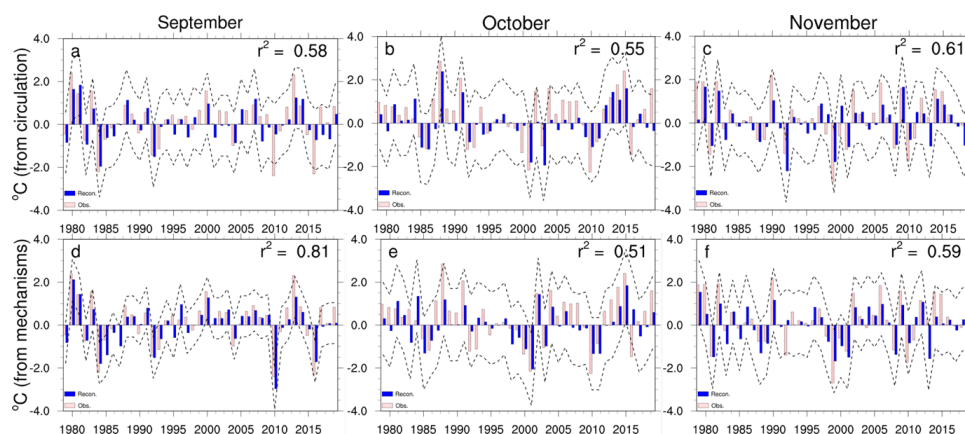


Figure 6. As in figure 4 (g-i), but using time series of key circulation features (south-west low, south-east low and southern Australian anticyclone) identified in figures 1 and 2 as predictors in the top row (a-c) and area-averaged dynamical heat mechanism components (850hPa zonal wind and meridional wind (multiplied by -1) and 500hPa vertical motion; see text for region averaged over) as predictors in the bottom row (d-e) for September, October, and November.

444

445 Following van Rensch et al (2019), indices of three dynamical heat mechanisms were
446 created by weighted area-averaging of westerly and northerly wind (meridional wind
447 multiplied by -1) over a region around southern Australia (25°S-45°S, 105°-155°E), and



448 500hPa vertical motion (ω ; positive is sinking motion) averaged over subtropical
449 Australia (15°S-25°S, 120°-155°E). Regions were selected based on the areas of highest
450 statistical significance between atmospheric circulation and Australian maximum
451 temperature in Fig. 2a-c. Again, a statistical model of Australian-averaged maximum
452 temperatures that uses these mechanisms as the predictors explains a higher proportion of
453 maximum temperature variance through spring than does the model using SAM or the
454 tropical TPI (Fig. 6 d-e). The percent variance explained is much higher in September (about
455 80%), before dropping to around 55% in October-November. The decrease in the percent
456 variance explained appears to be primarily associated with how strongly the westerly winds
457 correlate with maximum temperature over southern Australia; strong positive relationship
458 with westerly wind in September changes to insignificant or negative in October and
459 November (Supp. Fig S7 a-c). There is also an increase in negative correlation between
460 maximum temperatures and northerly winds in north-eastern Australia (Supp. Fig S7 d-e)
461 that will partly offset the increasing positive relationship further south. These changing
462 relationships between dynamical mechanisms and maximum temperature through spring
463 are linked with the changing relationships with the circulation features (Supp. Table 1)
464 through spring. Overall, however, the three dynamical heat mechanisms explain much of
465 Australia's monthly spring maximum temperature variability.

466

467 Figure 7 summarises the relationship between Australian maximum temperatures,
468 circulation features, dynamical heat mechanisms and climate drivers through the spring
469 months. The correlation between the SEC and Australian maximum temperature is
470 strongest in September and rapidly decreases through October and November, while
471 simultaneously the correlations with the SWC and particularly the SAA increase. As expected
472 from Fig. 2, the SEC and SWC are more closely linked with the extratropics. Linearly
473 regressing out the SAM component from time series of the SWC and SEC reduces the
474 correlation strength with Australian maximum temperature (Fig. 7a), particularly in
475 September. Conversely, linearly removing the tropical TPI slightly increases the correlation
476 between the cyclones and temperature, with the partial-correlation only weakening in
477 November. As SAM is strongly related to the barotropic cyclones it is also strongly related to
478 how temperature changes with the westerly wind. Linearly removing SAM from the
479 westerly wind time series nearly halves the correlation with maximum temperature in



480 September, and weakens the correlation in October and November (Fig. 7b). Conversely,
481 linearly removing the tropical TPI actually increases the correlation slightly with the westerly
482 wind in September and October, but decreases the correlation in November.
483
484

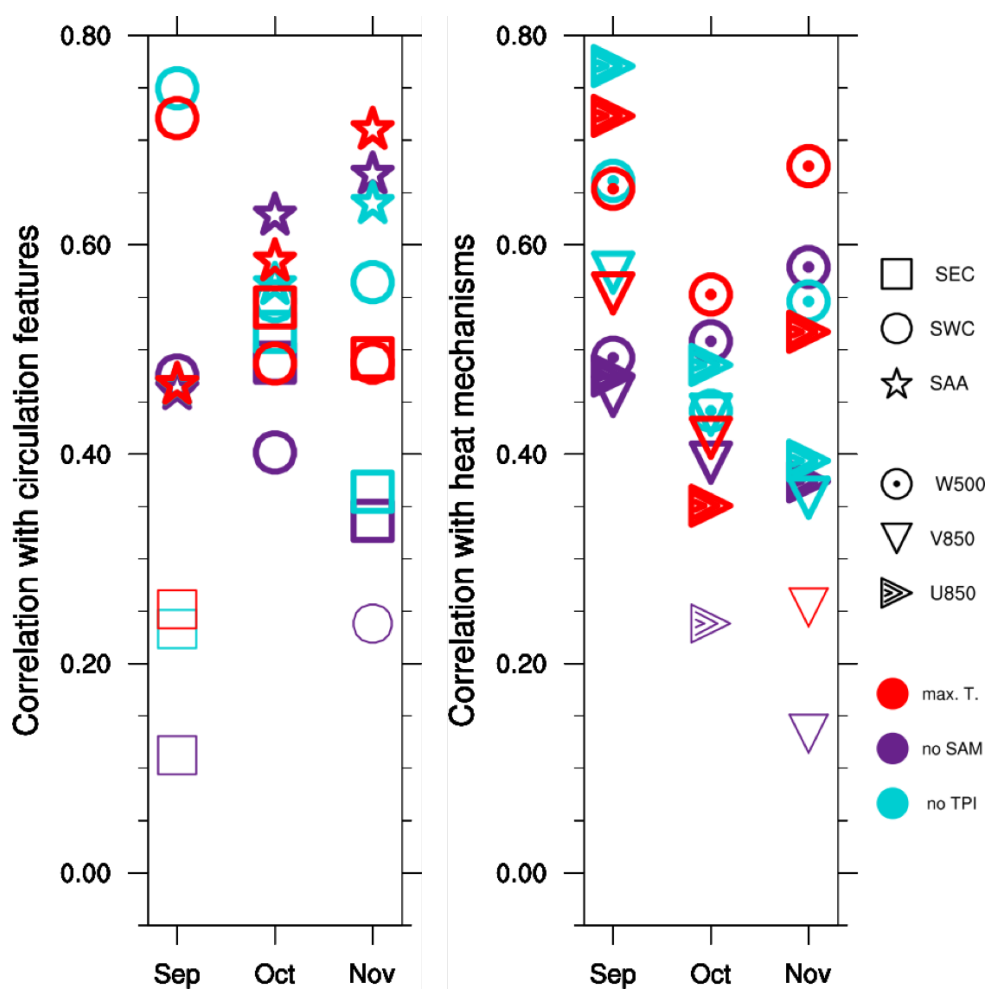


Figure 7. Correlations between Australian area-averaged maximum temperature (red) between key atmospheric circulation features (left figure) and dynamical heat mechanisms (right figure) for September, October and November. The purple and turquoise show partial correlations of the same, but with SAM and the tropical TPI linearly removed. Bold lines show the correlation was statistically significant at the 95% confidence level using a Student's t-test with 39 samples.



485 The relationships with northerly wind and sinking motion and Australian-averaged
486 maximum temperature do not change as dramatically with the removal of SAM or the TPI.
487 Northerly wind is not strongly influenced by the tropics or extratropics in September or
488 October, but the correlation strengthens and weakens in November with the removal of the
489 TPI and SAM, respectively. While removing SAM and TPI from the SAA had relatively little
490 influence on the correlation with Australian maximum temperatures, removing SAM from
491 sinking motion in September and both TPI and SAM in October and November reduced the
492 correlation. Overall, it appears that the heat mechanisms associated with high maximum
493 temperatures in spring are influenced differently by the different influence of the
494 extratropics and tropics on the local atmospheric circulation features through spring.
495

496 7. Discussion and conclusions

497 The sources of the atmospheric circulation pattern associated with high monthly-maximum
498 temperatures in Australia appear to change from primarily extratropical in early spring to
499 tropical forcing in late spring. Examination of three dynamical heat mechanisms (low-level
500 winds broken into westerly and northerly components, and mid-tropospheric sinking
501 motion) indicates that this shift may be due to a change in how heat develops. In early
502 spring, the low-level wind plays a greater role in maximum temperatures, advecting
503 relatively warmer air from the oceans over the cold land-mass. This wind correlates strongly
504 with the extratropics (here, SAM) as SAM projects strongly onto the southwest and
505 southeast cyclones that direct a lot of the low-level flow around Australia. Conversely, the
506 atmospheric circulation associated with the TPI ($x-1$) acts to counter the low-level flow that
507 drives higher temperatures. Thus, in early spring we have a closer association with heat
508 production and the extratropics. By late spring, the circulation patterns associated with high
509 temperature have changed and the wind does not correlate as strongly. As such adiabatic
510 sinking over subtropical Australia has a proportionally stronger correlation with high
511 temperatures. Both SAM and TPI ($x-1$) regressions show sinking motion in the subtropics
512 through spring, but it is the TPI that better matches the sinking motion over eastern
513 Australia in November. Hence, the apparent change from extratropical to tropical forcing in
514 the circulation pattern is because the tropics promotes more of the heat developing



515 mechanisms later in spring. However, much of the atmospheric patterns associated with
516 heat through spring are explained by neither the tropical TPI nor SAM,
517
518 The subtropical jet appears to play a greater role in Australian spring heat by acting as a
519 wave guide (Hoskins and Ambrizzi, 1993) that directs quasi-stationary Rossby waves toward
520 Australia, rather than as a block that limits direct propagation of Rossby waves from the
521 tropical Indian Ocean to the southern hemisphere extratropics (e.g. Simpkins et al., 2014; Li
522 et al., 2015 a,b). While wave activity flux only appears to propagate directly out of the
523 tropical Indian Ocean later in spring, this analysis does not suggest that the tropical Indian
524 Ocean is not a wave source in early spring. Indeed, the results are broadly consistent with
525 IOD-forced wave trains identified in the literature (Cai et al., 2011; McIntosh and Hendon,
526 2018; Wang et al., 2019). In particular, the secondary wave source in the high latitudes of
527 the Indian Ocean proposed by McIntosh and Hendon (2018) may be key for promoting the
528 TPI-forced atmospheric circulation in early spring, though this is beyond the scope of this
529 study to confirm. As the subtropical jet did not act as a barrier preventing the tropical Indian
530 Ocean's influence on Australia's maximum temperature, we argue instead that the apparent
531 change in forcing through spring was more related to the origins of three of the dynamical
532 heat mechanisms behind that heat. Consistent with this idea, wave activity flux calculated
533 by first regressing 200Z onto the three dynamical heat mechanisms (Sup. Fig. 8) also has
534 changing extratropical or tropical forcing through spring, that then propagates along the jet
535 wave guide toward Australia.

536
537 Area-averaged low-level wind and vertical motion were used to understand how the
538 atmospheric circulation relates to Australia-wide maximum temperatures, but do not form a
539 complete picture of spring temperature development in Australia. Statistical models using
540 these mechanisms explain much, but not all, of the maximum temperature variance over
541 Australia. Further, it was not always clear how the atmospheric circulation features
542 influenced those heat mechanisms. In particular, the southern Australian anticyclone and
543 500hPa subtropical-Australian sinking motion, while important for heat, appear to be largely
544 uncorrelated with the other circulation features and mechanisms. Greater insight into how
545 remote forcing of the atmospheric circulation results in high Australian temperatures could
546 be gained by including other heat mechanisms in future analyses, including: insolation (Lim



547 et al., 2019b), land-surface feedbacks linked to antecedent moisture (e.g. Arblaster et al.,
548 2014; Hirsch and King, 2020), and changes to synoptic weather systems (Cai et al., 2011;
549 Hauser et al., 2020). How each of these mechanisms relates to the others, and geographic
550 changes across Australia should also be considered. The combination of poleward advection
551 of adiabatically warmed air after it descended anticyclonically over the Tasman Sea has
552 been identified as a key mechanism for summer heatwaves in southeast Australia (e.g.
553 Quinting and Reeder, 2017). This combination of mechanisms may generate heat through
554 spring, particularly in the east and in November. The connection with rising motion over
555 southern Australia has also not been examined, and may indicate the importance of air
556 being diabatically warmed in association with storminess just to Australia's south, before
557 advecting and descending toward Australia. While the three heat dynamical heat
558 mechanisms were simple, the complex relationships between all of the mechanisms meant
559 that the three used in this analysis were broadly representative of a large portion of how
560 heat develops through spring.

561

562 We used the TPI to represent tropical variability relevant to Australia's maximum
563 temperature, but other indices or drivers may highlight different Rossby wave pathways or
564 heat mechanisms. Reconstructing Australian maximum temperature time series with more
565 commonly used indices for the IOD and ENSO did not change the effectiveness of the
566 statistical models overall (Supp. Fig. 4). However, it did suggest that the IOD had greater
567 influence on Australia's maximum temperature in early spring than does ENSO, consistent
568 with the seasonal-length studies of (Jones and Trewin, 2000; Saji et al., 2005). As such, we
569 may expect different monthly Rossby wave pathways to Australia associate with the IOD in
570 early spring, giving greater influence from the tropical Indian Ocean at this time. The MJO
571 generates Rossby wave trains from the western Pacific that promote low minimum
572 temperatures in Australia during winter (Wang and Hendon, 2020) and from the tropical
573 Indian Ocean to promote high maximum temperatures in Australia in spring (personal
574 communication: Wang and Hendon, 2021). The positive phase of the IOD suppresses MJO
575 activity across the Indian Ocean (Wilson et al., 2013), possibly restricting the MJO's
576 influence on Australia's maximum temperature at such times. However, MJO activity in the
577 tropical Indian Ocean has recently been found to counter the wetting influence of La Niña
578 during spring (Lim et al., 2021b). As such the MJO may be an important factor for spring



579 maximum temperatures when the tropical SSTs are not otherwise conducive for high
580 temperatures, but is beyond the scope of this study.

581

582 As the trend toward higher Australian spring temperatures is projected to continue into the
583 future a better understanding of what drives maximum temperatures over the months of
584 spring is critical for better prediction and better preparation to adapt to a warming climate.
585 A combination of extreme values in remote drivers of variability, including extreme positive
586 IOD, central-Pacific El Niño, and sustained negative SAM associated with very strong sudden
587 stratospheric warming, exacerbated already dry and hot conditions in spring 2019 to
588 promote one of Australia's deadliest fire seasons (Watterson, 2020; Lim et al., 2021a;
589 Abram et al., 2021, Marshall et al. 2021). Further, projected trends toward positive IOD (Cai
590 et al., 2014; Abram et al., 2020) or toward negative TPI (Timbal and Hendon, 2011) may
591 contribute to higher maximum temperatures in the future, particularly in later spring when
592 the tropics exert greater influence on Australia's dynamical heat mechanisms. As we have
593 shown just how different the atmospheric circulation and heat mechanisms can be through
594 a season in Australia, other regions and seasons could also benefit from similar analysis,
595 particular as the world continues to warm (e.g. Collins, et al., 2013).

596

597 Code and data availability

598 The code for analysis is available from the corresponding author on request. ERA5-
599 reanalysis data are available from Copernicus Climate Change Service at
600 <https://www.ecmwf.int/en/forecasts/datasets/reanalysis-datasets/era5>. AWAP data is
601 available from the Australian Bureau of Meteorology.

602

603 Author contribution

604 R.M.C produced the figures and wrote the initial draft manuscript. All authors contributed
605 to analysis and editing of the manuscript.

606



607 Competing interests

608 The authors declare that there are no conflicts of interests.

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625

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