

1 **A critical evaluation of decadal solar cycle imprints in the MiKlip**
2 **historical ensemble simulations**

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9

10 **Abstract**

11 Studies concerning solar-terrestrial connections over the last decades claim to have found evidence that the quasi-
12 decadal solar cycle can have an influence on the dynamics in the middle atmosphere in the Northern Hemisphere
13 during winter season. It has been argued that feedbacks between the intensity of the UV part of the solar spectrum and
14 low latitude stratospheric ozone may produce anomalies in meridional temperature gradients which have the potential
15 to alter the zonal-mean flow in mid to high latitudes. Interactions between the zonal wind and planetary waves can
16 lead to a downward propagation of the anomalies, produced in the middle atmosphere, down to the troposphere. More
17 recently, it has been proposed that top-down initiated decadal solar signals might modulate surface climate and
18 synchronize the North Atlantic Oscillation. A realistic representation of the solar cycle in climate models was
19 suggested to significantly enhance decadal prediction skill. These conclusions have been debated controversial since
20 then due to the lack of missing realistic decadal prediction model set ups and more extensive analysis.

21 In this paper we aim for an objective and improved evaluation of possible solar imprints from the middle atmosphere
22 to the surface and with that from head to toe. Thus, we analyze model output from historical ensemble simulations
23 conducted with the state-of-the-art Earth system model MPI-ESM-HR. The target of these simulations was to isolate
24 the most crucial model physics to foster basic research on decadal climate prediction and to develop an operational
25 ensemble decadal prediction system within the MiKlip framework.

26 Based on correlations and multiple linear regression analysis we show that the MPI-ESM-HR simulates a realistic,
27 statistically significant and robust shortwave heating rate and temperature response at the tropical stratopause, in good
28 agreement with existing studies. However, the dynamical response to this initial radiative signal in the NH during the
29 boreal winter season is weak. We find a slight strengthening of the polar vortex in midwinter during solar maximum
30 conditions in the ensemble mean, which is consistent with the so-called “top-down” mechanism. The individual
31 ensemble members, however, show a large spread in the dynamical response with opposite signs in response to the
32 solar cycle, which might be a result of the large overall internal variability compensating rather small solar imprints.

33 We also analyze the possible surface responses to the 11-year solar cycle and review the proposed synchronization
34 between the solar forcing and the North Atlantic Oscillation. We find that the simulated westerly wind anomalies in
35 the lower troposphere as well as the anomalies in the mean sea level pressure are most likely independent from the
36 timing of the solar signal in the middle atmosphere and the alleged top-down influences. The pattern rather reflects

37 the decadal internal variability of the troposphere, mimicking positive and negative phases of the Arctic- and North
38 Atlantic Oscillations throughout the year sporadically, which is then assigned to the solar predictor time series without
39 any physical plausible connection and sound solar contribution.

40 Finally, by applying lead/lag correlations, we find that the proposed synchronization between the solar cycle and the
41 decadal component of the North Atlantic Oscillation might rather be a statistical artefact, affected for example by the
42 internal decadal variability of the ocean, than a plausible physical connection between the UV solar forcing and quasi-
43 decadal variations in the troposphere.

44

45 1. Introduction

46 The discipline of decadal climate prediction is rather young and a rapidly growing field in climate science. By using
47 initialized climate model simulations, the gap between weather forecasting and long-term climate model projections
48 covering the complete 21st century or beyond is bridged (e.g., Pohlmann et al., 2013; Meehl et al., 2014). By the aid
49 of decadal climate predictions, policymakers can be equipped with an improved decision-making basis allowing for a
50 better planning of necessary water resources, agriculture, energy and infrastructure measures for the near-term future
51 (Mehta et al., 2011). The aim of the German joint research project “Mittelfristige Klimaprognose” (MiKlip) was to
52 establish a new decadal prediction system allowing for a more precise midterm climate forecasting. To this effect,
53 potential driving factors shaping the decadal climate from both anthropogenic and natural sources have been evaluated
54 critically based on large ensemble simulations with the Max Planck Institute for Meteorology Earth System Model
55 (MPI-ESM).

56 One factor that potentially influences tropospheric weather and climate is the variability in the middle atmosphere via
57 stratosphere-troposphere coupling processes. The internal variability in the middle atmosphere during the dynamically
58 active winter and spring seasons is strongly controlled by the variability of Rossby waves, which propagate upward
59 from the troposphere to the middle atmosphere where they break and interact with the zonal-mean flow. The changes
60 in the zonal-mean flow, again, can alter the propagation conditions for planetary scale waves initiating a self-consistent
61 feedback called wave-mean flow interaction (e.g. Andrews 1985). As a result, strong disruptions, born in the middle
62 atmosphere, such as sudden stratospheric warmings (SSWs), which are characterized by a breakdown of the polar
63 vortex, have the potential to propagate downward into lower atmospheric layers and interfere with the tropospheric
64 weather regime (e.g., Baldwin and Dunkerton, 2001). A prominent example for this are Northern Hemisphere (NH)
65 cold air outbreaks which have the tendency to be more frequent and severe in seasons with a weak stratospheric polar
66 vortex (e.g. Huang et al., 2021).

67 A source of variability that might influence the dynamics in the middle atmosphere on the decadal timescale via a
68 complex feedback mechanism between radiation, chemistry and wave-mean flow interaction is the 11-year solar cycle.
69 Pioneering work concerning the impact of the solar cycle on middle atmosphere dynamics and possible connections
70 to the troposphere goes back to Kodera and Kuroda (2002). Based on a relatively short period of NCEP reanalysis data
71 (1979 – 1998), the authors observed an increase of the tropical stratopause temperature (TST) (at ~50 km) during the

72 time of the solar maximum. In their conceptual explanation, this temperature increase leads to a strengthening of the
73 meridional temperature gradient and an intensification of the polar night jet (PNJ) in the winter stratosphere. The
74 stronger westerlies create a barrier for upward propagating planetary waves, which in turn are deflected poleward and
75 break at lower altitudes. The resulting divergence in the Eliassen-Palm flux (EPF) allows the positive wind anomaly
76 to move downward and poleward over the winter season. Kodera (2002) argues that the solar induced wind anomalies
77 may advance into the troposphere, where they create a signal in meteorological variables mimicking a positive phase
78 of the North Atlantic Oscillation (NAO). Matthes et al. (2004, 2006) studied the proposed “top-down” mechanism by
79 the aid of idealized simulations with an early 3-dimensional middle atmosphere general circulation model (GCM).
80 Analysing monthly to sub-monthly means, they found that during solar maximum conditions the polar vortex seems
81 to be stronger especially in November and December and linked this to a positive Arctic oscillation (AO)-like pattern
82 which they found in lower altitudes and to some extent at the surface. The observed pattern weakens in January and
83 changes sign from February onwards. In subsequent studies comparable results have been found (e.g., Schmidt et al.,
84 2010; Ineson et al., 2011; Chiodo et al., 2012; Langematz et al., 2013). However, the exact timing of the progression
85 of the signals from the middle atmosphere to the surface depends on the individual study and varies from December
86 to February. A most recent study again concludes that the most pronounced solar signal seem to appear in early winter
87 (Ma et al. (2018)). These early studies are often quoted as convincing proof for a “top-down” influence of the 11-year
88 solar cycle in both the middle atmosphere and the troposphere. Complementary to this, Gray et al. (2013) found that
89 the strongest NAO-like solar-induced signals in the North Atlantic (i.e. a positive phase of the NAO) actually seem to
90 appear with a time lag of three to four years after the solar maximum in the respective seasonal winter mean (DJF).
91 However, the observed lags could not be reproduced in coupled atmosphere-ocean simulations conducted by the same
92 group. In the model, the postulated response to the solar cycle in the North Atlantic appears almost in phase with the
93 solar forcing (maximum response between lag year zero to one) (Gray et al., 2013). This discrepancy between observed
94 and simulated lag in the response in the North Atlantic NAO was confirmed in subsequent studies (e.g., Scaife et al.,
95 2013; Andrews et al., 2015).

96 With respect to possible solar induced impacts on NH surface variability in the winter season, Thiéblemont et al.
97 (2015) went one step further. Analyzing a simulation incorporating 150 model years, they claim that the solar forcing
98 synchronizes the decadal component of the NAO variability spectrum, a phase relation they cannot find in an
99 experiment without 11-year solar variability. This result has been debated controversially since its publication. Chiodo

100 et al. (2019) found almost identical spectra of the NAO decadal variability in two simulations of 500 model years each,
101 with and without a 11-year solar cycle forcing. Furthermore, they identified NAO patterns in similar time segments in
102 both experiments (forced and unforced). They suspect, therefore, that the alleged surface solar signals in other studies
103 are most likely a result of the internal variability of the NAO itself rather than solar cycle imprints. On the other hand,
104 Drews et al. (2022) most recently argue that the solar cycle near-surface imprints can only shine through during very
105 active solar periods with large amplitudes of the 11-year solar cycle. They also state that during these periods the
106 surface decadal prediction skill would be significantly enhanced if the solar cycle is a vital part of the prediction
107 system. In the context of the most recent literature, it is difficult to understand why Chiodo et al. (2019) and Drews et
108 al. (2022) arrive at a different assessment of the solar signal, even though the same model was used. This might point
109 to the fact, that the complexity of the model is not the most relevant component in shaping potential surface solar
110 signals, but rather the effects of internal variability in individual model runs and (to some degree) the applied analysis.

111 In this publication, we evaluate possible imprints of the 11-year solar cycle in different domains of the atmosphere
112 from the initial solar radiative signal in the tropical upper stratosphere down to the surface in the NH winter season.
113 We analyze the MiKlip historical ensemble simulations conducted with the state-of-the-art Earth system model MPI-
114 ESM-HR, which is the physical basis for the decadal prediction system, which is operational at the “Deutscher
115 Wetterdienst” (DWD) since 2020. The availability of the large amount of output data from the MiKlip historical model
116 ensemble enables us to address the unresolved questions of the solar surface imprint, such as the dependence of the
117 signal on the solar cycle amplitude, on a more robust statistical basis than is possible in single model simulations. In
118 our study, we aim to identify the role of the solar imprints for the decadal variability of the NAO in winter. While the
119 model simulations include both, changes in the total solar irradiance (TSI) and spectral solar irradiance (SSI), potential
120 effects related to solar energetic particles (SEP) and medium energy electrons (MEE) are not explicitly included in the
121 MiKlip experiments. Observations and model studies suggest that changes in the stratospheric composition related to
122 SEP can lead to a radiatively driven modulation of the middle atmosphere dynamics, which can penetrate to lower
123 atmospheric layers down to the troposphere (e.g., Seppälä et al., 2009, 2014; Baumgaertner et al., 2010; Arsenovic et
124 al., 2016). However, since no robust surface impacts have been simulated even for strong solar energetic particle
125 events (SEP) of the recent decades (Jackman et al., 2009), we infer that including these effects may not alter our results
126 significantly.

127 This publication is structured as follows. In Section 2 we describe the MPI-ESM, the setup of the analyzed simulations
128 and the applied methodologies to detect potential solar cycle signals in different atmospheric domains. In Section 3,
129 the initial radiative solar signal in the tropical middle atmosphere is evaluated. Subsequently, we concentrate on the
130 dynamical response to the initial solar signal in the NH winter season. Here we show in Section 4 the ensemble mean
131 response and compare individual ensemble members with opposite solar signatures. In Section 5, we derive solar-
132 induced signals near the surface in our simulations and observations. In Section 6, we check our model results with
133 respect to the proposed synchronization between the solar forcing and the NAO. Finally, we summarize and discuss
134 our results in a broader context (Section 7).

135

136 2. Data and methods

137 2.1 Model description and experimental design

138 The historical simulations analyzed in this publication have been conducted with the Max Planck Institute for
139 Meteorology Earth System Model in high resolution configuration (MPI-ESM1.2-HR; hereafter called MPI-ESM-
140 HR) at the Deutsches Klimarechenzentrum (DKRZ). MPI-ESM-HR includes the atmospheric general circulation
141 model ECHAM (European Centre Hamburg) version 6.3 (ECHAM6.3) with a horizontal/vertical resolution of
142 T127L95 (corresponds to a ~ 100 km * 100 km model grid and 95 levels in the vertical with a model top at 0.01 hPa
143 or ~ 80 km) (Müller et al., 2018). The high vertical resolution allows for an internally generated quasi-biennial
144 oscillation (QBO) in the tropical stratosphere (Pohlmann et al., 2019). Radiative processes are represented using the
145 rapid radiation transfer model for GCMs (RRTM-G) for both the shortwave and longwave part of the electromagnetic
146 spectrum (Iacono et al., 2008). Other diabatic processes, such as vertical mixing by turbulence and moist convection,
147 large-scale convection, and momentum deposition by orographic and unresolved gravity waves are described in more
148 detail in Stevens et al. (2013). Oceanic processes are accounted for in the coupled Max Planck Institute ocean model
149 (MPIOM) with a TP0.4 (0.4° nominal) resolution (Jungclaus et al., 2013). MPI-ESM-HR further incorporates the
150 biogeochemistry module Hamburg Model of the Ocean Carbon Cycle (HAMOCC) (Ilyina et al., 2013; Paulsen et al.,
151 2017) and the land surface model JSBACH (Reick et al., 2013).

152 In this publication, we analyze 10 members of the MPI-ESM-HR historical simulations performed within the German
153 research project MiKlip. The MiKlip historical ensemble simulations include the observed natural and anthropogenic

154 climate drivers, as described in the CMIP5 protocol (Taylor et al., 2013). The individual ensemble members (1 to 10)
155 have been initialized from different model years of a 1850 preindustrial (PI) control simulation and were integrated
156 over the period 1850 to 2005. Here, we focus on the period 1880 – 1999. Thus, a total of 1,200 model years have been
157 evaluated. Since the model does not include interactive atmospheric chemistry, ozone concentrations have to be
158 prescribed. In the MiKlip historical simulations, the merged CMIP5 ozone dataset was used, which consists of a
159 combination of SAGE I+II satellite and radiosonde data in the period 1979 to 2005. To derive earlier ozone
160 concentrations back to 1850, the zonal mean stratospheric time series is extended backwards based on the regression
161 fits and proxy time series of equivalent effective stratospheric chlorine (EESC) and solar variability (Cionni et al.,
162 2011). The solar variability forcing includes all observed solar cycles and follows Lean (2000).

163

164

165 2.1 Data analysis

166 *Detrending, correlations, filtering*

167 To detrend the sunspot number (SSN) (Source: WDC-SILSO, Royal Observatory of Belgium, Brussels -
168 <https://www.sidc.be/silso/infosnmtot>) and shortwave heating rate time series, a third-degree polynomial function has
169 been fitted to the data, the respective anomalies are shown in Figure 1 (the original, unfiltered SSN time series is
170 shown in Supplementary Figure 1). The detrended SSN time series has then been correlated (Pearson r) with the
171 detrended tropical stratopause temperature (defined as the mean value between 25°S – 25°N at 1 hPa (Figure 3). All
172 correlation analyses have been performed by using the Python `scipy.pearsonr` function. Statistical significance of the
173 correlations has been calculated by using a two-tailed Student's t-test, as implemented in Python. To reduce the
174 degree of internal variability, a Butterworth bandpass filter with cutoff frequencies of 9 and 13 years has been
175 applied to the detrended PNJ time series (defined as the arithmetic mean of the zonal-mean zonal wind between
176 35°N – 45°N at 1 hPa) (Figure 3). The same Butterworth bandpass filter has also been applied to the zonal-mean
177 zonal wind time series at 10 hPa (zonal mean over 55°N – 65°N) (Figure 3) and the NAO time series. The NAO time
178 series has been calculated by the aid of an EOF analysis conducted for the MSLP data over the Atlantic sector (20 –
179 80°N, 90°W – 40°E) in the winter season (DJF averaged and individually for December, January and February). The
180 first principal component is then used to describe the NAO variability. The lead/lag correlations (Figure 8) are then
181 calculated between the filtered NAO and SSN time series.

182 *Multiple linear regression*

183 To detect the solar cycle signals in the middle atmosphere (Figures 2, 4 and 5) and in the mean sea level pressure in
184 both observations and model data (Figures 6 and 7), we use an established multiple linear regression (MLR) technique
185 as described in Bodeker et al. (1998). To derive the individual regression coefficients, we use a set of six predictors in
186 the MLR model:

187
$$X(t) = \text{Off.const} + A * \text{CO2}(t) + B * \text{QBO}(t) + C * \text{QBOorth}(t) + D * \text{SSN}(t) + E * \text{Nino3.4}(t) + F * \text{tau}(t) + R(t)$$

188 with: Off.const = annual cycle; CO2(t) = increase in the atmospheric CO₂ concentration; QBO(t) = phase of the QBO,
189 defined by the zonal-mean zonal wind in 30 hPa (5°S – 5°N); QBOorth(t) = the orthogonal of QBO(t); SSN(t) = SSN
190 time series; Nino3.4(t) = Nino3.4 times series; tau(t) = optical thickness at 550 nm and R(t) = model residuum. Based
191 on this MLR analysis, we derived the model response to our chosen set of predictors, e.g., the temperature response
192 per unit of the predictor (i.e., K per 1 SSN). To display the model response during solar maximum, we scaled the
193 coefficients to 180 SSN, which is a good approximation for a mean solar cycle amplitude between 1880 and 1999. To
194 detect potential time lags in the response to the solar cycle at the surface, the solar time series has been shifted in such
195 a way that the model response lags the solar forcing by 1 to 4 years.

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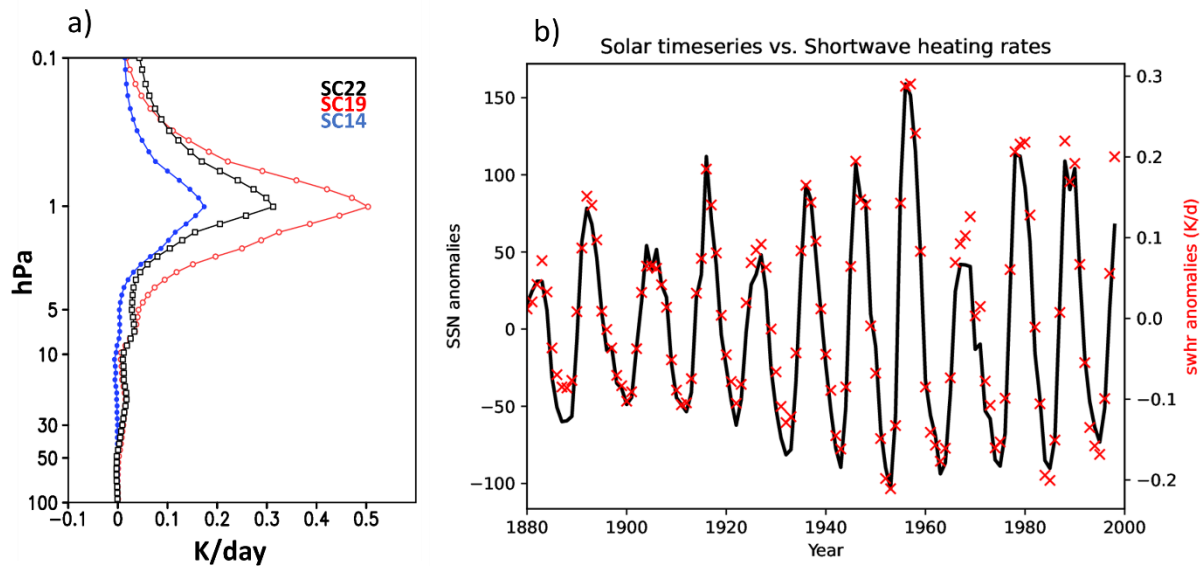
197 3. The initial radiative solar signal in MPI-ESM

198 The dynamical “top-down” mechanism, assumed to be the pathway for the propagation of the solar signature through
199 the atmosphere to the surface in NH winter (see also Section 1), is initiated at the tropical upper stratosphere by the
200 absorption of solar ultraviolet (UV) irradiance by ozone and molecular oxygen. In particular, the absorption of solar
201 photons by ozone in the Hartley bands (200 – 310 nm) in the upper stratosphere - and to a lesser extent the Huggins-
202 bands (310 nm – 400 nm) in the middle stratosphere – heats the upper stratosphere increasingly with height and leads
203 to the formation of the warm stratopause. Although the variation in solar UV-irradiance over the 11-year solar cycle
204 is less than 10% in the ozone absorption bands, the enhanced UV radiation at solar maximum – in combination with
205 increased ozone concentrations - leads to stronger shortwave heating and a concurrent warming of the tropical
206 stratopause by the order of 1 K, as has been derived from merged MSU4 and SSU+MLS-satellite observations (Randel
207 et al., 2016).

208 Figure 1a shows the annual mean response of the modelled shortwave radiative heating rate (SWHR) at the
209 stratosphere and lower mesosphere (100 – 0.1 hPa) for a range of solar cycle (SC) amplitudes from the weak SC14 (in
210 blue), over the medium SC22 which has been used as solar forcing in the CMIP5 protocol (in green), to the very strong
211 SC19 (in red). MPI-EMS-HR produces the well-known solar cycle impact with enhanced SW heating during solar
212 maximum throughout the upper stratosphere and lower mesosphere. The maximum SWHR difference develops at the
213 stratopause and ranges for the three selected solar cycles between 0.17 and 0.51 K/day. With a SWHR increase of 0.32
214 K/day for the SC22 solar forcing, MPI-ESM-HR produces an initial solar radiative response at the tropical stratopause

215 which is in very good agreement with offline radiation model calculations using the CMIP5 solar forcing (i.e., the
 216 same forcing as in MPI-ESM-HR) in a line-by-line reference and two CCM (EMAC and WACCM) radiation codes
 217 (see Figure 8, yellow curves in Matthes et al., 2017). This is a significant improvement compared to the earlier
 218 ECHAM4 and ECHAM5 model versions which were not able to simulate the SWHR response to the solar cycle in the
 219 stratosphere (see Figure 17 in Forster et al., 2011), and thus missed the initial solar temperature signal necessary for
 220 the “top-down” mechanism. The improvement in the MPI-ESM-HR is the result of the enhanced spectral resolution
 221 of the new shortwave radiation scheme in ECHAM6 which resolves the shortwave spectrum in 14 bands spanning the
 222 wavelength range from 820 to 50,000 cm^{-1} (Iacono et al., 2008), whereas ECHAM4 and ECHAM5 used a lower
 223 spectral resolution with the four-band model of Fouquart and Bonnel (1980), later extended to six bands by Cagnazzo
 224 et al. (2007).

225 Figure 1b shows the time series of the SSN and the modeled SWHR at the tropical stratopause over the period from
 226 1880 – 1999. The shown anomalies of both time series from a third-degree polynomial fit clearly demonstrate that
 227 solar cycles of different amplitudes initiate SWHR responses that closely follow in magnitude the strength of the solar
 228 forcing. Only during SC20, the maximum SWHR response is higher than expected for that weak solar cycle. This is
 229 not reproduced in the SWHR, possibly due to the transition from synthetic SSN before 1979 to observed SSN
 230 afterwards.



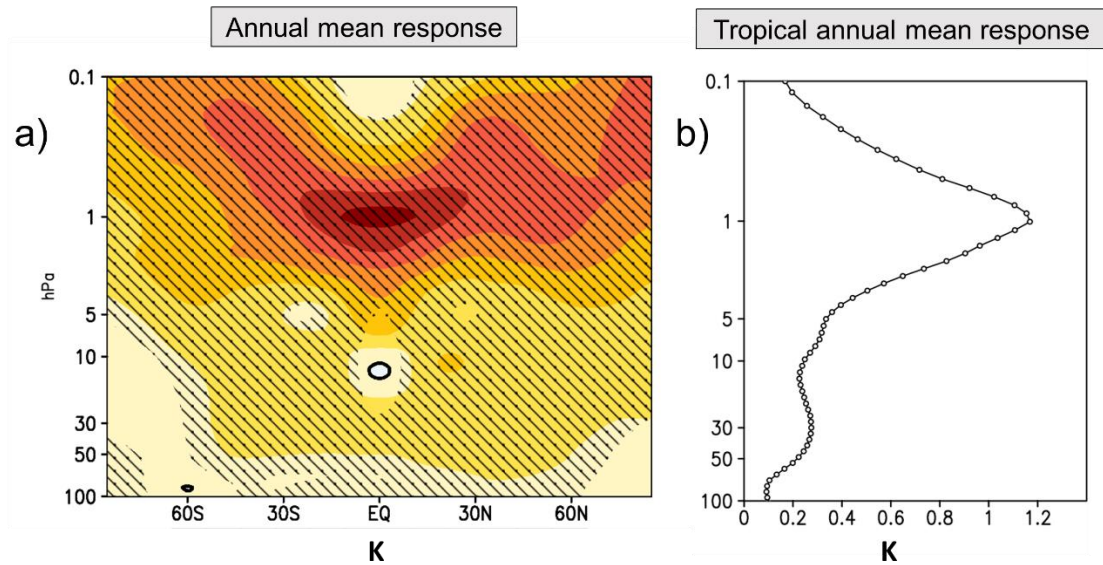
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232 **Figure 1:** Solar shortwave heating rate signature in the MPI-ESM-HR historical simulations: a) Annual tropical
 233 mean (25°S – 25°N) shortwave heating rate difference in K/day between the maximum and minimum of three solar

234 cycles: the weak solar cycle 14 (blue), the medium solar cycle 22 used in CMIP5 (green), and the strong solar cycle
 235 19 (red) (a), and: Time series of the sunspot number and the annual tropical mean (25°S – 25°N) shortwave heating
 236 rate at the stratopause (1 hPa). Shown are anomalies from a third-degree polynomial fit to the data (b).
 237

238 When averaging over all solar cycles between 1880 and 1999 and all 10 ensemble members, we obtain a robust, highly
 239 significant annual mean warming of the complete middle atmosphere at solar maximum (Figure 2a), reaching a peak
 240 response of 1.2 K at the tropical stratopause (Figure 2b). This result is slightly higher than the solar signal derived
 241 from satellite observations (0.7 K / 100 solar flux units), respectively ~1 K between solar minimum and maximum)
 242 (Randel et al., 2016). Given the excellent temporal evolution of the initial radiative response of the upper tropical
 243 stratosphere to the decadal solar forcing, we conclude that MPI-ESM-HR produces the necessary prerequisite for the
 244 dynamically enhanced “top-down”-mechanism, which will be investigated in more detail in the next section.

245
 246



247 **Figure 2:** Long-term annual ensemble mean response based on MLR analysis of the zonal-mean temperature (in K)
 248 to the solar cycle in the middle atmosphere as a function of height and latitude (hatched regions mark the 95% level
 249 of significance) (a), and the annual mean tropical (25°S – 25°N) temperature response (in K).
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253 4. Downward transfer of the solar signal to the surface: the key role of dynamics

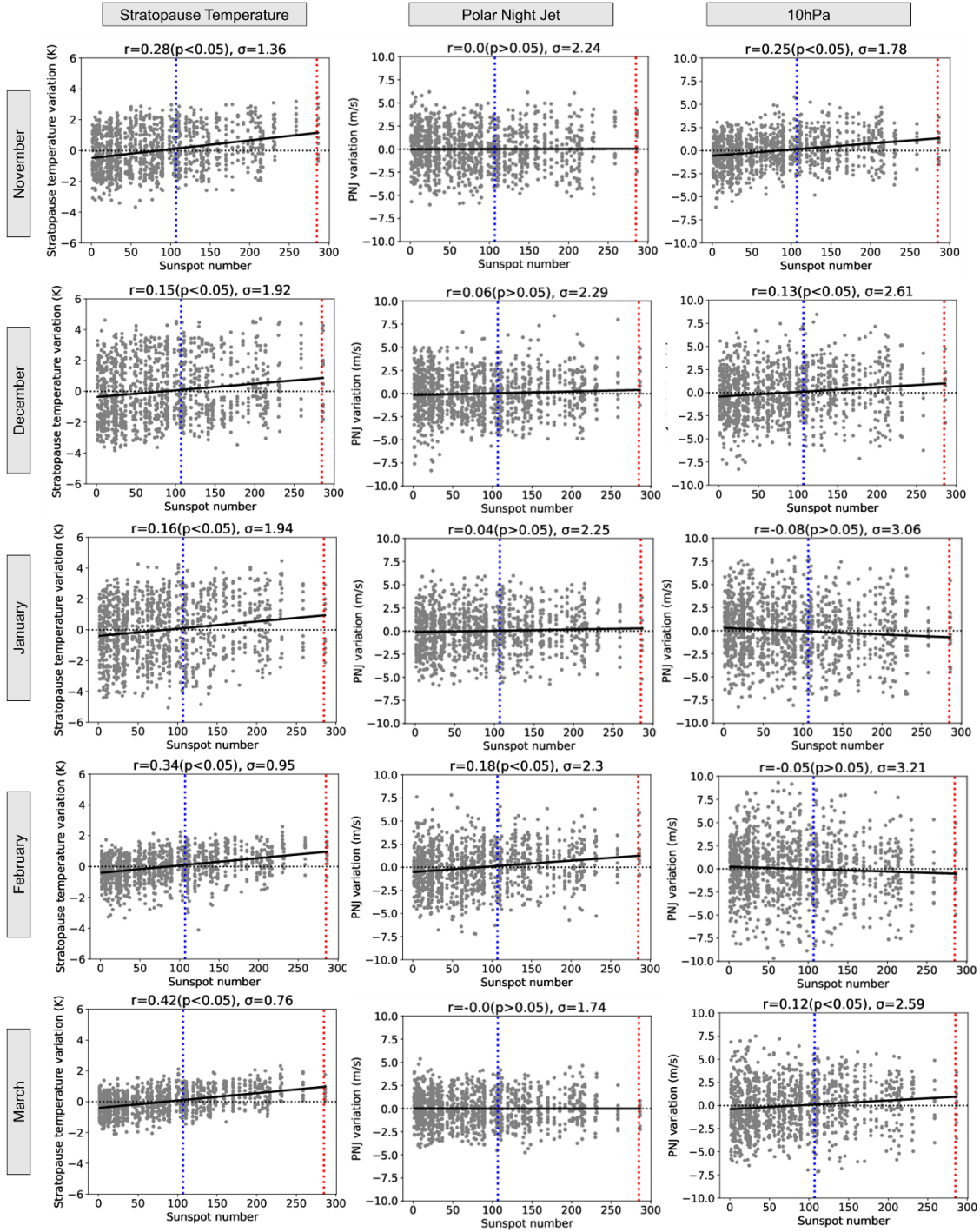
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255 After having demonstrated the ability of the MPI-ESM-HR model to realistically simulate the radiative and the related
 256 temperature response in the tropical upper stratosphere to the decadal solar forcing, we investigate as next step the
 257 potential dynamical reaction to the radiative forcing, which is expected according to the “top-down” mechanism. By
 258 evaluating the ensemble spread in the NH during the dynamically active season (November to March), we assess the

259 variability of different dynamical variables in the stratosphere with respect to the solar fluctuations in the MPI-ESM-
260 HR historical ensemble simulations. We focus first on the detrended deviations from the long-term monthly means for
261 the TST and (to estimate the dynamical response in the NH) the zonal-mean zonal wind at two different altitudes and
262 latitudes (Figure 3). To approximate the PNJ (the local maximum wind speed in the upper stratosphere) we use the
263 mean of the zonal-mean zonal wind in 35°– 45°N at 1 hPa. The variability in the middle stratosphere is represented
264 by the mean of the zonal-mean zonal wind in 55°– 65°N at 10 hPa. After calculating the respective anomaly time
265 series for the TST, the PNJ and the 10 hPa zonal wind variations for each month individually, we correlate these time
266 series with the detrended DJF mean SSN time series. To mute the interannual variability (operating on timescales
267 between 1 and 8 years) of the polar vortex, the PNJ and 10 hPa anomaly time series, as well as the SSN time series,
268 have been bandpass-filtered, before calculating the correlations. Our results indicate that the TST correlates
269 significantly with the SSN, not only in the annual mean (compare Figure 1b) but also in each individual month
270 considered (Figure 3, left column). While negative and positive TST anomalies (i.e., negative and positive deviations
271 from the long-term monthly mean) are almost uniformly distributed for SSN values smaller than the SC14 maximum
272 (blue dotted lines), an increase in the solar forcing exceeding the SC14 SSN maximum leads to a higher probability of
273 positive TST anomalies. The strength of the correlations changes over the season, such that a stronger connection
274 between the solar forcing and the temperature response at the tropical stratopause is given in late autumn (November:
275 $r=0.28$) and late winter (February: $r=0.34$; March: $r=0.42$). In these months, a particular strong solar forcing (indicated
276 by the SSN value of the SC19 maximum (red dotted lines)) is almost always associated with a positive temperature
277 anomaly at the tropical stratopause. Weaker correlations and a broader distribution of negative and positive
278 temperature anomalies, even during periods with especially pronounced solar activity, are calculated for the midwinter
279 season (December: $r=0.15$; January: $r=0.16$). These findings are consistent with an increase in the overall variability
280 in the TST during December and January, making it more difficult for the relatively weak solar induced signals to be
281 distinguished from the background noise. The higher variability in the TST during December and January is probably
282 a result of the higher variability of the tropical branch of the Brewer-Dobson circulation (BDC) in boreal winter (e.g.,
283 Butchart, 2014).

284 According to the general concept of the “top-down” mechanism the initial signal in the TST would be accompanied
285 by a strengthening of the PNJ via a modification of the meridional temperature gradients. Considering the statistically
286 significant temperature signals and correlations at the tropical stratopause in the MPI-ESM-HR model (Figure 3, left

287 column), we expect a dynamical response of the PNJ in our simulations. However, the correlations between the SSN
288 and the PNJ time series (Figure 3, middle column) do not show statistically meaningful relations between the solar
289 forcing and the dynamical response of the PNJ. Only during February, a weak but statistically significant correlation
290 is found, which might be related to the enhanced impact of the solar forcing in the TST during the same month.
291 However, this connection as well becomes insignificant, if the correlations are calculated based on the unfiltered SSN
292 and PNJ time series. Figure 3 (right column) shows the correlations between the solar forcing and the zonal mean
293 zonal wind for the lower (and more northward) 10 hPa anomaly time series. We find the strongest (and significant)
294 correlations in November ($r=0.25$) and December ($r=0.13$), although these correlations become (again) negligible if
295 the correlations are calculated based on unfiltered model data. The differences in the timing between the maximum
296 correlations of the SSN with the PNJ (February) and the 10 hPa zonal wind time series (November and December) are
297 not in line with the established idea of a successive “poleward and downward” progression of the dynamical solar
298 signal. Furthermore, the computed SSN/PNJ correlations for November, December, January and March are ≤ 0.06 ,
299 implying that the characteristics of the PNJ are not markedly influenced by the magnitude of the solar forcing and thus
300 the amplitude of the solar cycle.



301

302 **Figure 3:** Scatter diagram of the stratopause temperature (left column), PNJ (middle column) and zonal-mean zonal
 303 wind averaged over $55^{\circ}\text{N} - 65^{\circ}\text{N}$ at 10 hPa (right column) variations vs. SSN. The numbers given in the headings
 304 show the correlation coefficients (r), their statistical significance ($p < 0.05$: significant correlation, or $p > 0.05$:
 305 insignificant correlation), and the overall variation (σ). The dotted blue and red lines indicate the SSN at solar cycle
 306 maximum for SC14 and SC19 (the weakest/strongest solar cycles considered in the simulations).
 307
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309 Figure 3 demonstrates that while the connection between the solar forcing and the TST is clearly visible in our
310 correlation analysis, the potential dynamical response in the NH is harder to detect, especially due to the highly variable
311 polar vortex. Therefore, we proceed using a MLR analysis to separate the potential dynamical solar induced signals
312 from other internal generated disturbances in the ensemble mean.

313 After having analyzed the variability of the TST, the PNJ and the 10 hPa zonal-mean zonal wind, we will now isolate
314 potential solar signals by the aid of MLR. Figure 4 shows the solar regression coefficients, scaled to a mean amplitude
315 of the solar cycle (180 SSN), for the zonal-mean temperature (top row), the zonal-mean zonal wind (middle row) and
316 the EPF (vectors) and its divergence EPFD (colors) (bottom row) for each NH winter month (November – March).

317 Here, we focus on the potential solar cycle signals between the equator and the North Pole and pressure heights in
318 1.000 hPa – 0.1 hPa for the temperature and wind responses and 100 hPa – 0.1 hPa for the EPF diagnostics. We find
319 a significant response in the zonal mean temperature at the tropical stratopause (Figure 4, top row) with a maximum
320 response at the equator of 1.2 K during November. The solar induced temperature signal is confined to the inner tropics

321 in late autumn and early winter and advances towards higher latitudes between January and March. This is consistent
322 with the seasonal march of the incidence angle of solar radiation after the winter solstice in December. In the middle
323 to polar latitudes, we find a clear dipole in the temperature anomalies especially during November and December.

324 This dipole is characterized by distinct (and significant) positive temperature anomalies in the lower mesosphere and
325 upper stratosphere and weak (and insignificant) negative anomalies in the middle and lower stratosphere. Particularly
326 the pronounced polar heating in the upper stratosphere from November to December agrees well with a most recent
327 analysis of ERA-interim reanalysis data by Kuroda et al., (2022). The detected temperature signals in the middle

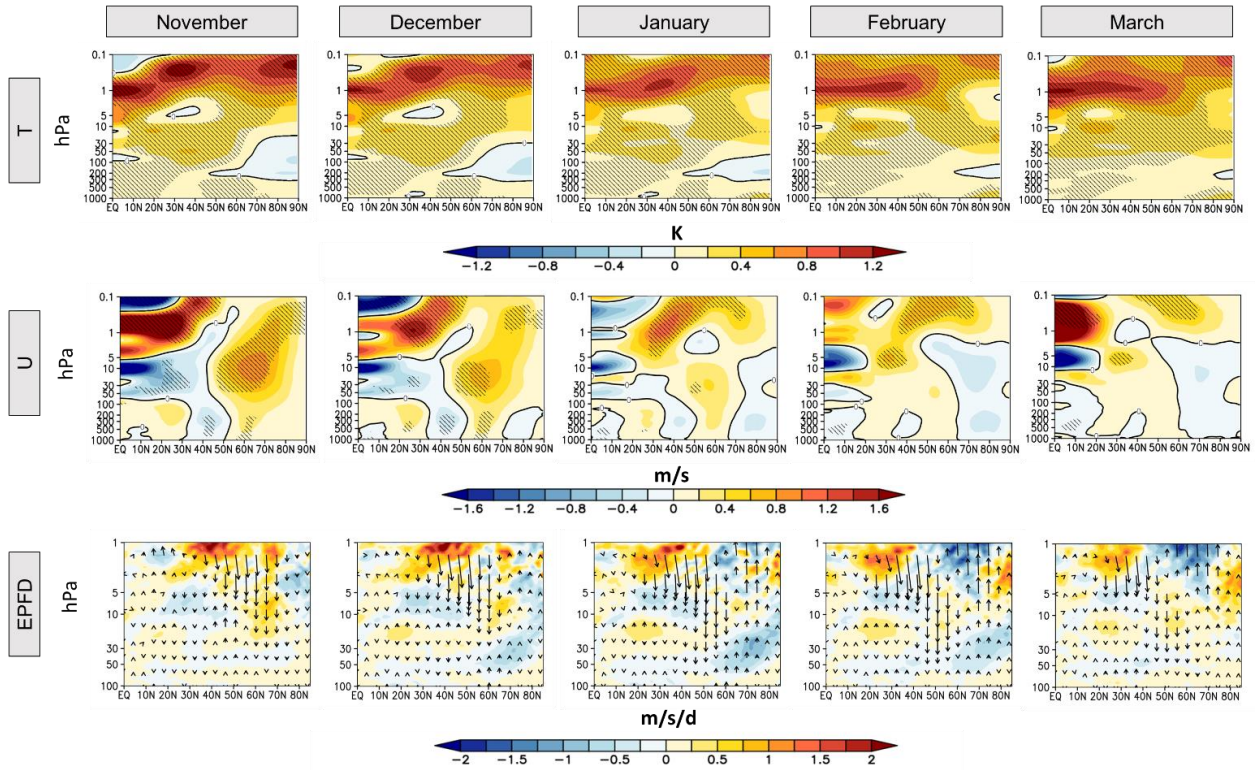
328 atmosphere in November and December are in line with the anomalies in the zonal-mean zonal wind (Figure 4, middle
329 row), which indicate a stronger (and thus cooler) polar vortex during these months. Additionally, a convergence of the
330 EPF (indicated by the reddish colors in Figure 4, bottom row) and its (here downward oriented) vectors imply a reduced
331 upward propagation of planetary waves due to the strengthening of the polar vortex. The maximum (and significant)

332 response in the stratospheric zonal-mean zonal wind in the area of the polar vortex, is located at $\sim 60^\circ\text{N}$ at 10 hPa.
333 Here, we find positive anomalies of the zonal-mean zonal wind of ~ 1 m/s. Given the mean zonal-wind speeds between
334 20 m/s (November) and 30 m/s (December), simulated by the model (not shown) at this height and latitude, the solar
335 influence seems rather small in comparison. The detected dipole in the zonal-mean temperature starts to weaken from

336 January on and vanishes almost completely until March. During the same months, we find a (yet insignificant)

337 weakening of the polar vortex which allows for more upward propagation of planetary waves (indicated by a
338 divergence of the EPF (bluish colors) and upward oriented vectors). In the troposphere, a weak (≤ 0.5 m/s) but
339 significant westerly wind anomaly around $\sim 60^\circ\text{N}$ can be detected in November and December. The weak tropospheric
340 wind response agrees with other studies (Matthes et al., 2006; Schmidt et al., 2010; Ineson et al., 2011; Chiodo et al.,
341 2012; Langematz et al., 2013; Kuroda et al., 2022; Drews et al., 2022).

342 While in some studies the march of the westerly wind anomalies from the middle atmosphere to the surface seems to
343 follow the proposed “poleward and downward” concept (e.g., Matthes et al., 2006; Ineson et al., 2011; Drews et al.,
344 2022), the signal transmission in the MPI-ESM-HR and other model simulations (e.g., Schmidt et al., 2010; Chiodo et
345 al., 2012; Kuroda et al., 2022) rather follows a “downward-only” storyline. Additionally, the description of the
346 westerly wind anomalies at the surface is sometimes inconsistent with the idea of a successive downward propagation
347 of the signal from higher to lower altitudes. As an example, significant westerly wind anomalies at the surface at
348 middle latitudes are already present in November in the modeling studies of Matthes et al. (2006) and Kuroda et al.
349 (2022), even though the major signal is still high up in the middle atmosphere. Furthermore, in Kuroda et al. (2022)
350 the westerly wind anomalies at the surface at middle latitudes are present throughout the complete season (i.e., in all
351 months between November-March), similar to our MPI-ESM-HR simulations. In other studies, the westerly anomalies
352 are insignificant (e.g., Schmidt et al., 2010) or do not reach the ground (e.g., Chiodo et al., 2012). This implies that the
353 detected surface wind anomalies could be independent from the seasonal march in the middle atmosphere and might
354 rather be a product of the internal variability in the troposphere (i.e., the AO or NAO) itself. Likewise, the temperature
355 response to the solar cycle in the troposphere with positive temperature anomalies of ≤ 0.2 K at the surface is rather
356 weak (Figure 4, top row). Interestingly, these small temperature signals are significant in the tropics in all considered
357 months, which is consistent with the high (and relatively constant) solar insolation in the inner tropics and a damped
358 overall variability compared to the extratropical regions. By contrast, the significant surface temperature anomalies in
359 the extratropical regions are located between 50°N and 60°N until January and shift towards the polar latitudes in
360 February and March.



361

362 **Figure 4:** The ensemble mean long-term response (based on MLR) to the solar cycle of the zonal-mean temperature
 363 (first row), zonal-mean zonal wind (second row) (hatched regions mark the 95% level of significance), and the EPF
 364 (vectors) and the divergence of the EPF (EPFD, colors) in the NH during the boreal winter season. All results have
 365 been scaled to 180 SSN.

366

367

368 So far, we focussed on the discussion of the potential solar signals in the ensemble mean derived from the 10 individual
 369 MiKlip historical simulations thus obtaining statistically more robust results than is possible through analyses of single
 370 simulations. The necessity of working with ensemble mean results is impressively demonstrated by comparing two of
 371 our 10 individual ensemble members. Figure 5 shows the solar regression coefficients for the zonal-mean temperature
 372 and zonal-mean zonal wind for the ensemble members 1 (EM1, top panel) and 4 (EM4, bottom panel), as in Figure 4.
 373 The derived patterns for the solar zonal-mean temperature signal in EM1 show distinct similarities with the ensemble
 374 mean. As an example, we find a (significant) maximum temperature response around the tropical stratopause.
 375 Furthermore, the distribution of the temperature anomalies in the middle to higher latitudes again displays the polar
 376 heating in the lower mesosphere and the upper stratosphere and the cooling in the middle to lower stratosphere. Again,
 377 this pattern starts to weaken from January on. We notice that in comparison to the ensemble mean, fewer areas depict
 378 significant temperature signals, even though the magnitude of the temperature response is stronger. This can be

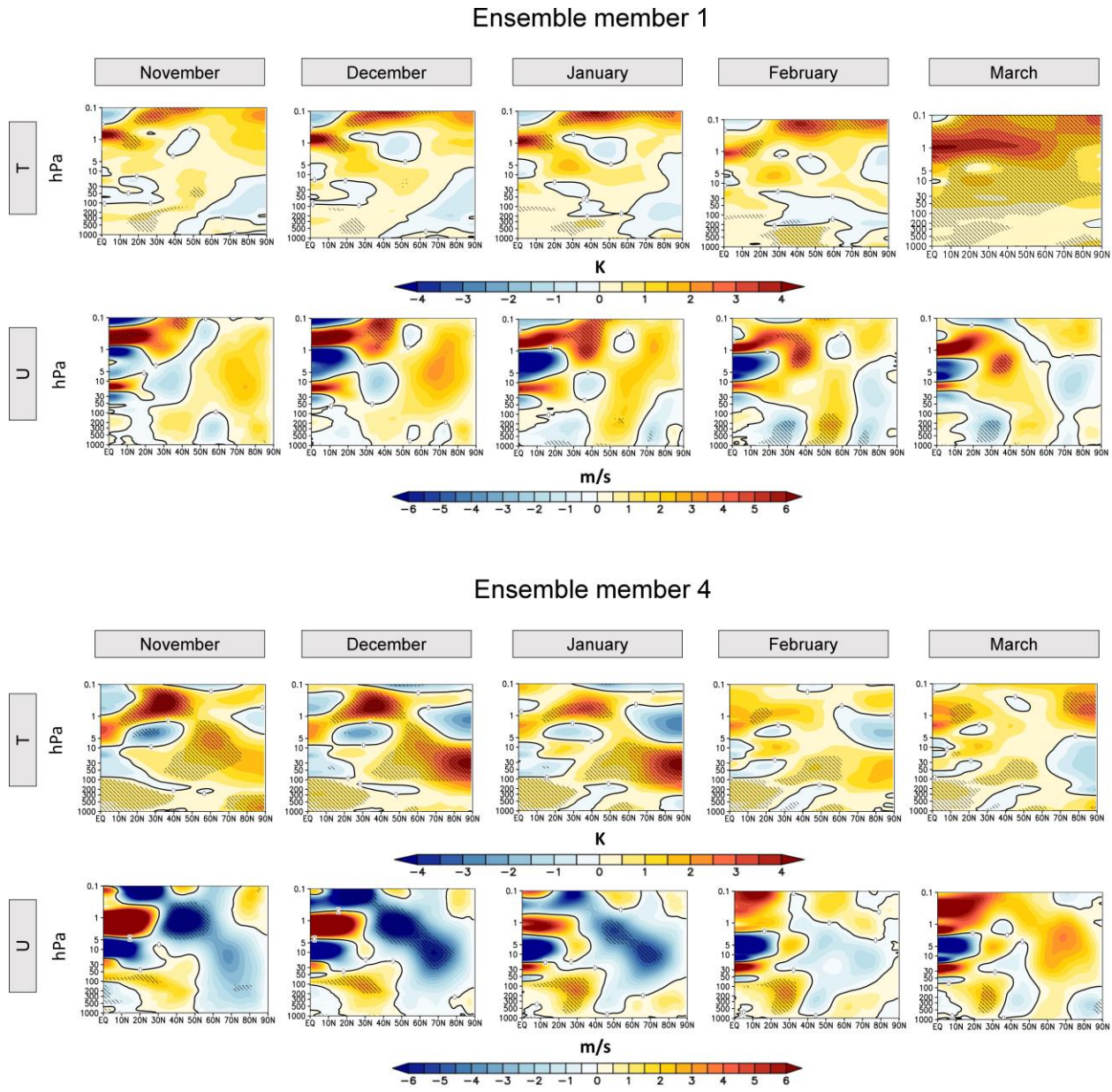
379 attributed to the fact that the analysis only includes 120 model years and thus ~12 solar cycles (instead of 1.200 and
380 ~120 in the ensemble mean), which is seemingly not enough to dampen the internal variability and inhibits the solar
381 induced signals to become significant against the overall background noise. Likewise, the solar response of the zonal-
382 mean zonal wind in the middle atmosphere in EM1 shows the main characteristics, as already noticed in the ensemble
383 mean, such as a strengthening of the polar vortex in November and December and a subsequent weakening and a
384 conversion in sign afterwards. However, none of the detected signals in the area of the polar vortex are statistically
385 significant. As for the response of the zonal-mean zonal wind at the surface, we detect significant anomalies in January
386 and February. The geographical distribution of the anomalies (westerly wind anomalies at middle latitudes and easterly
387 wind anomalies at polar latitudes), however, mimic a negative phase of the AO which is not in line with the general
388 concept of solar induced “top-down” influences.

389 In EM4, the initial temperature signal in the upper tropical stratosphere is, as in EM1, visible throughout the complete
390 season and the strongest in November and December. Thus, the response to the solar cycle in these latitudes and
391 heights turns out to be a robust feature in the MPI-ESM-HR model experiments. However, even though exactly the
392 same solar forcing has been applied in EM4 as in EM1, the dynamical response of EM4 looks very different. For
393 instance, we find a cooling of the polar upper stratosphere and a (significant) warming in the middle to lower
394 stratosphere in December and January. This pattern is common during SSWs, which (by chance) could have been
395 more frequent in EM4 during December and January than in EM1. The strong and significant easterly wind anomalies
396 in the middle atmosphere, indicating a slowdown of the polar vortex during these months, underpin this hypothesis.
397 These findings imply that the detected signals in EM1 could also be a result of (by chance) less frequent SSWs in EM1
398 leading to a potentially misleading attribution to solar variability. In our simulations, four out of 10 simulations show
399 a weakening of the polar vortex during high solar activity, while six depict a strengthening of the latter, which may
400 explain the rather weak tendency to westerly wind anomalies in the ensemble mean.

401 Either way, our results point to the fact that the internal dynamics of the polar vortex have the ability to control the
402 transmission of potential solar induced signals from the tropics to the polar regions and are thus more important than
403 the amplitudes of individual solar cycles (compare also Figure 3), as recently claimed by Drews et al. (2022).

404

405



406

407 **Figure 5:** Long-term response (based on MLR) to the solar cycle of the zonal-mean temperature (first row) and the
408 zonal-mean zonal wind (second row) (hatched regions mark the 95% level of significance) in the two ensemble
409 members EM1 (top panels) and EM4 (bottom panels) in the NH during the boreal winter season. All results have been
410 scaled to 180 SSN.

411

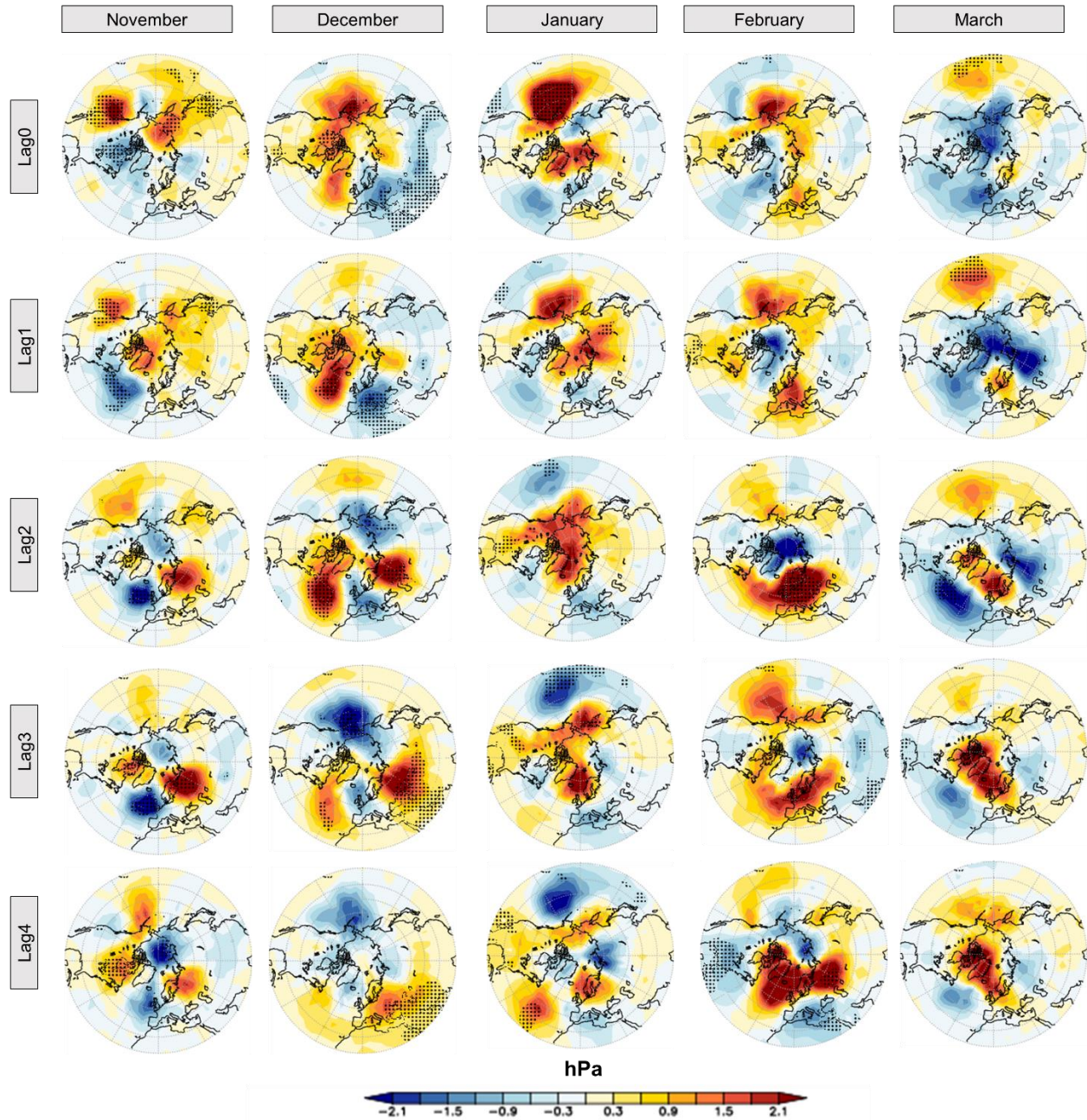
412

413 **5. Direct and lagged surface solar signals**

414 Our results so far indicate a robust response of the TST to the quasi-decadal solar cycle. The subsequent dynamical
415 response in the NH during the boreal winter season, however, is difficult to assess. By the aid of a MLR analysis we
416 could detect weak solar cycle imprints in the zonal-mean temperature and the zonal-mean zonal wind in the ensemble
417 mean. However, these signals are not robust among all individual ensemble members, especially with respect to the
418 detected anomalies in the zonal-mean zonal wind at the surface which seem to be independent of the signals in the
419 middle atmosphere.

420 Nevertheless, in the next step, we first aim at detecting potential solar signals at the surface by applying the MLR
421 analysis to mean sea level pressure (MSLP) data in NH winter. Figure 6 shows the monthly solar regression
422 coefficients for MSLP, scaled to a mean solar cycle amplitude of 180 SSN, in the HadSLP2 observational dataset
423 (Allan and Ansell, 2006) for the same period as simulated (1880 – 1999). In order to check for eventual time lags
424 between the applied solar forcing and the model response, as suggested for example by Gray et al. (2013), lagged
425 regressions were calculated by shifting the solar predictor time series against the observations so that it leads the model
426 data between one and four years. Our results show positive and negative anomalies in the MSLP in the middle and
427 polar latitudes which mimic positive and negative phases of the AO in a rather random than systematic way. As an
428 example, we find an AO-positive like pattern (i.e., negative pressure anomalies over the North Pole and positive
429 pressure anomalies in the surrounding middle latitudes) in November at lag year four, in December at lag year four,
430 in February at the lag years one to three and in March at lag year one. The most pronounced AO-positive anomalies,
431 with a negative but insignificant anomaly of ~ 2 hPa over the North Pole and a positive anomaly of the same magnitude
432 in the middle latitudes, are given at lag year 2. Hence, the strength of the detected potential solar signals in our
433 HadSLP2 analysis is in line with other studies assessing observational products (e.g., Gray et al., 2013; Kuroda et al.,
434 2022; Drews et al., 2022). The detected maximum impact at lag year 2 in February in our analysis, however, agrees
435 with Kuroda et al. (2022) and Drews et al. (2022) but differs from Gray et al. (2013) who found a maximum response
436 at lag year 4 in the DJF mean. These discrepancies in the timing of the peak solar-induced surface signal in the HadSLP
437 MSLP data can only be explained by differences in the analysis techniques, and reveal a high sensitivity of solar-
438 induced surface signals to the applied methodology and individual interpretation of the results.

HadSLP2

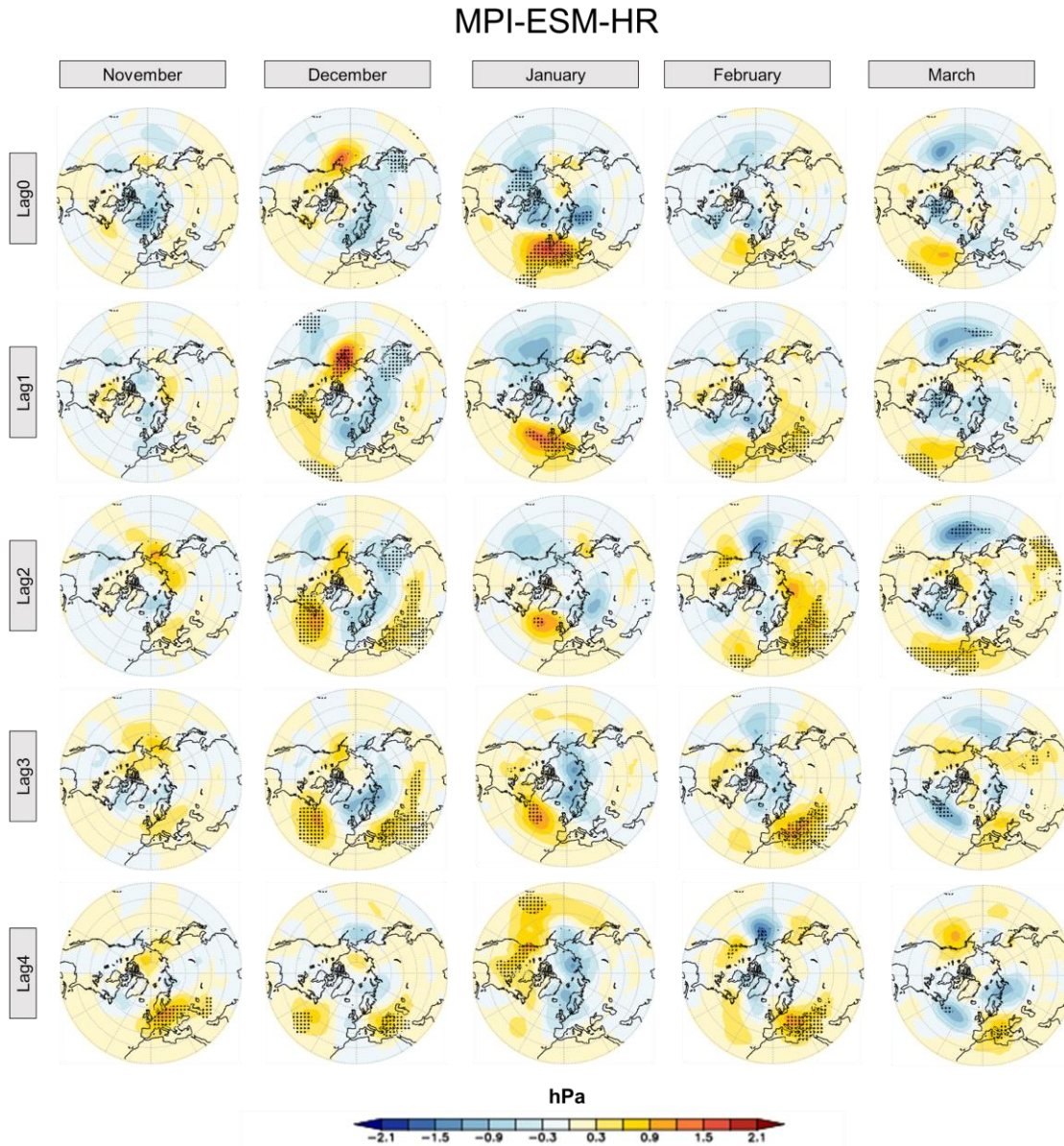


439

440 **Figure 6:** The (lagged) response of mean sea level pressure (MSLP) to the solar cycle in the NH during the boreal
441 winter season for the HadSLP2 dataset (dotted regions mark the 95% level of significance). Columns denote the
442 individual months of the winter season; rows indicate the lag of the MSLP time series with respect to the solar forcing
443 time series.

444

445 Figure 7 shows the same analysis for the MiKlip historical simulations, i.e., the ensemble mean of the solar regression
446 coefficients for the MSLP for each month (November to March) and the (lag) years zero to four. We detect AO-
447 positive-like anomalies in the MSLP in December at the lag years 0 and 1, in January at the lag years 0 to 4 and in
448 February at the lag years zero to four. The strongest negative MSLP anomalies over the North Pole show a response
449 of ~ -1.5 hPa and $\sim +1.5$ hPa in the middle latitudes in January and December. Thus, the overall model response is
450 weaker compared to the observational data. This is not surprising given the fact that the model results depict the mean
451 over 10 ensemble members (with respective dampening effects) compared to one ‘ensemble member’ representing the
452 observations. While the detected magnitudes of the MSLP anomalies in MPI-ESM-HR agree with other solar cycle
453 model studies (e.g. Gray et al., 2013; Scaife et al., 2013; Andrews et al., 2015; Drews et al., 2022), the detected timing
454 (i.e. the progression of the signals from the middle atmosphere to the surface) in the MPI-ESM-HR does not fit the
455 narrative of the “top-down” mechanism as described most recently by Kuroda et al. (2022) and Drews et al. (2022).
456 In these studies, the authors find the most pronounced AO-positive like pattern in February at the surface and link this
457 to the coupling between the stratosphere and the troposphere, which peaks in exactly this month. In contrast, in our
458 model simulations the strongest coupling between the stratosphere and the troposphere appears in December (see
459 Figure 4), while the most pronounced AO-positive like patterns appear in January and February at different lag years.
460 We, therefore, conclude that the detected surface solar signals could rather be a product of the internal variability in
461 the troposphere itself than being necessarily a consequence of the proposed “top-down” mechanism. Even if we assume
462 that the detected surface signals have a pure solar source (and the “top-down” mechanism is always present during
463 solar maximum years) it seems to be questionable in our view if these tiny signals would have the capability to
464 synchronize powerful large-scale climate modes such as the AO or the NAO, if they only emerge once per decade
465 over the duration of a month. As an example, the Icelandic Low and the Azores High, both controlling the pressure
466 gradients in the North Atlantic sector, show a month-by-month variation of ~ 8.5 hPa and ~ 6 hPa during winter time
467 in the model (not shown).



468

469 **Figure 7:** As Figure 6, but for the ensemble mean of the MPI-ESM-HR MiKlip historical simulations.

470

471 **6 A synchronization of the NAO by the solar cycle?**

472 In the following, we will address the question, if the quasi-decadal variations of the solar cycle have the ability to
 473 synchronize the decadal component of the NAO, as proposed by Thiéblemont et al., (2015) and Drews et al., (2022).

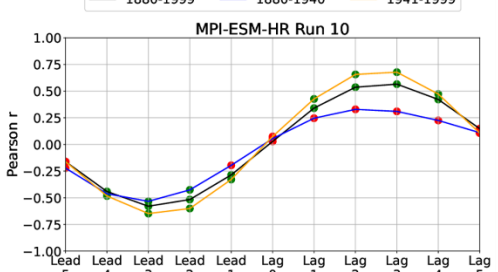
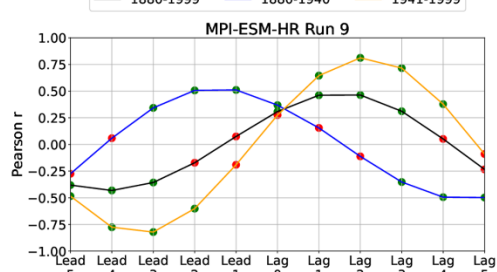
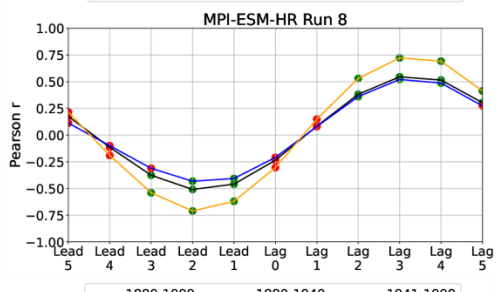
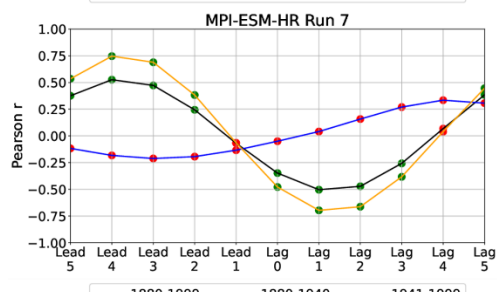
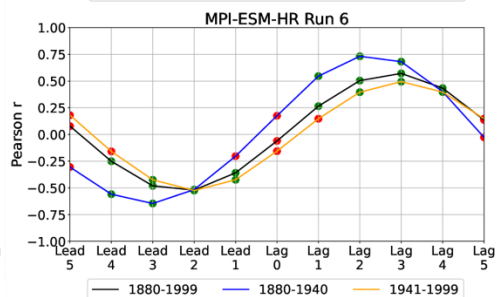
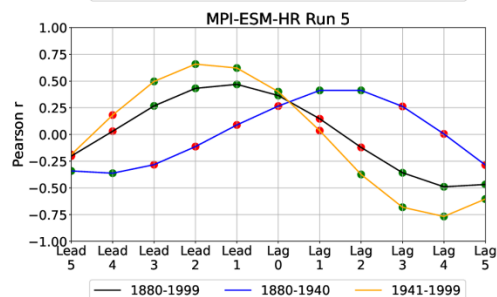
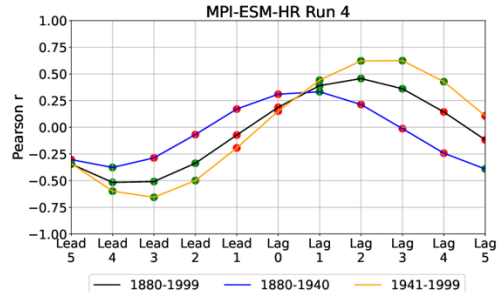
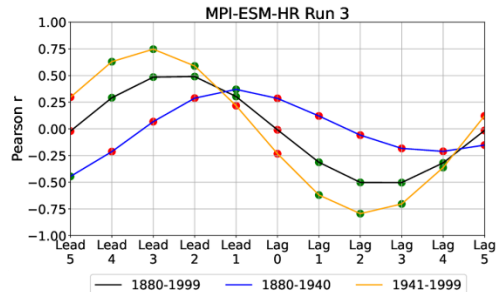
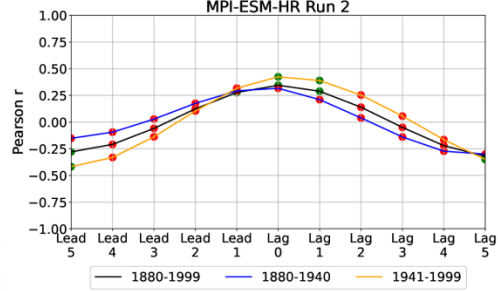
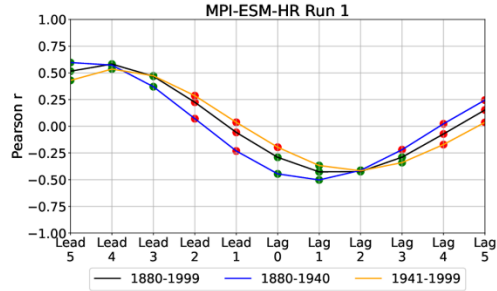
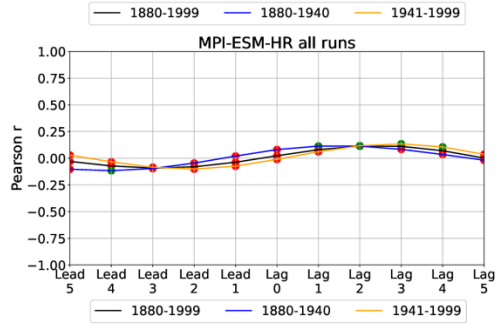
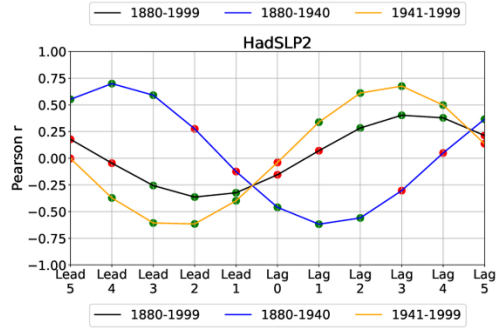
474 For a better comparison, we apply the same analytical strategy as proposed by Thiéblemont et al. (2015) to our model
 475 simulations and the HadSLP2 data, however with the exception that we use the SSN instead of the F10.7 solar flux

476 times series as a solar proxy. Since both the SSN and F10.7 time series show the same oscillations on the interannual
477 and decadal timescale, this is irrelevant for the interpretation of the results. First, an EOF analysis is applied to the
478 deseasonalized MSLP data over the Atlantic sector ($20 - 80^{\circ}\text{N}$, $90^{\circ}\text{W} - 40^{\circ}\text{E}$) in the winter season (DJF averaged).
479 The resulting leading principal components (PC1) are then used to describe the variability of the NAO. To mute major
480 parts of the interannual variability, we apply a Butterworth bandpass filter with cutoff frequencies of 9 and 13 years
481 to the PC1 and the SSN time series. As a result, the filtered PC1 and SSN time series only include the oscillations
482 operating on the quasi-decadal timescale. Subsequently, lead/lag correlations are calculated between the bandpass-
483 filtered PC1 and SSN timeseries for both the complete dataset and all individual ensemble members (1 to 10). Drews
484 et al. (2022) recently argued that the correlations would become more meaningful during the course of the 20th century
485 due to a series of solar cycles with stronger amplitudes. We, therefore, compute the correlations for three different
486 time segments: the whole period (WP) (1880 – 1999), the early period (EP) with weaker solar amplitudes (1880 –
487 1940) and the late period (LP) with more pronounced solar amplitudes (1941 – 1999).

488 For the HadSLP2 dataset (Figure 8, left column/first row) positive correlations between the decadal variation of the
489 NAO and the solar forcing is found for the lag years one to four in both the WP and the LP periods, with maximum
490 correlations at lag year three during the LP. For the EP, we find an out-of-phase relation between the solar time series
491 and the NAO on the decadal timescale. The evaluation of this (1 ensemble member) observational dataset implies that
492 the solar forcing actually leads the surface response by a couple of years and that this relation is more pronounced
493 during phases of higher solar activity. Indeed, similar phase relations in the different time segments are given in
494 individual ensemble members of the MiKlip historical simulations (e.g., EM9 (Figure 8, left column/sixth row).
495 However, phase relations like these seem far from being a robust feature if all model runs are considered. As an
496 example, EM5 (Figure 8, left column/third row) indicates positive correlations between the decadal behavior of the
497 SSN and the NAO time series for the lag years one to three during the EP, while this relation reverses (showing
498 negative correlations) during the WP and LP. This is also true for EM3 (left column/third row) and EM7 (left
499 column/fifth row). Other ensemble members (EM2; Figure 8, right column/second row) suggest a maximization of the
500 solar impact at lag year zero and this independently of the considered period. Furthermore, EM6 (Figure 8, right
501 column/fourth row) indicates stronger positive correlations at positive lag years during the EP than during the LP. The
502 most striking discrepancies, however, come from EM1 (Figure 8, left column/second row) and EM4 (Figure 8, right
503 column/third row). While EM1 shows negative correlations between the solar forcing and the NAO at positive lags

504 (in all time segments), this is vice versa in EM4. These surface responses in EM1 and EM4 are, however, opposite to
505 what would be expected from the polar vortex responses in these two ensemble members (a pronounced strengthening
506 of the polar vortex and a downward propagation of westerly wind anomalies to the surface in EM1, and a weakening
507 of the polar vortex and a downward propagation of easterly wind anomalies to the surface in EM4 during winter (see
508 Figure 5)) and opposite to the ‘top-down mechanism’.

509 When applied to the complete dataset of the MiKlip historical simulations, the correlation analysis yields a weak
510 positive (albeit significant) correlation at the lag years two to four, rather independently of the considered time
511 segment. This, however, should rather be interpreted as a slight (and by chance) overhang to positive correlations in
512 the MiKlip dataset (that could change in a larger ensemble) than a robust physical connection between the solar forcing
513 and the NAO. To verify whether the use of the seasonal mean (DJF) might dampen the solar cycle response, as
514 discussed by Drews et al. (2022), we repeated the analysis for the individual winter months (December, January and
515 February, see Supplementary Figure 21) for the model data. We did not detect stronger connections between the decadal
516 solar forcing and the NAO in the calculations based on individual months compared to the seasonal mean. On the
517 contrary, the correlation analysis based on the December months (i.e., the month where we find the “strongest” “top-
518 down” signals in the middle atmosphere) depicts negative correlations at positive lag years. In summary, given all of
519 these inconsistencies we suspect that there is no robust connection between the quasi-decadal solar oscillations and
520 the respective phase of the NAO in the CMIP5 MiKlip historical ensemble simulations.



522
523 **Figure 8:** Lead/lag-correlations between the seasonal mean (DJF) bandpass filtered PC1 based on NAO and SSN
524 time series. For the HadSLP2 dataset and the ensemble mean of the MPI-ESM-HR historical simulations (top row)
525 and the individual MPI-ESM-HR historical runs (rows 2 to 6) for different periods. Green dots mark statistically
526 significant (95%) correlations.
527

528

529 **7. Summary and discussion**

530 Our analysis of the MiKlip historical ensemble simulations, conducted with the state-of-the-art Earth system model
531 MPI-ESM-HR, revealed robust (and statistically significant) solar signals in the TST (see Figures 1 and 2). The
532 dynamical response to the initial solar temperature signal at the tropical stratopause, in the NH middle to polar latitudes
533 during the boreal winter season, however, showed a large spread among our data. This applies to the variability of the
534 PNJ and the 10 hPa zonal-mean zonal wind time series, which both did not show meaningful correlations with the
535 solar forcing (see Figure 3). When removing other than decadal variability components by MLR analysis, we were
536 able to detect (albeit rather weak) solar signals in the NH winter, in both the ensemble mean zonal-mean temperature
537 and zonal-mean zonal wind, that basically agree with the proposed “top-down” influence of solar variability in the
538 middle atmosphere (see Figure 4). However, the MLR analysis based on individual ensemble members revealed
539 signals of opposite direction (i.e., a strengthening (EM1) or weakening (EM4) of the polar vortex during periods of
540 high solar activity) (see Figure 5). Furthermore, we find indications that the detected anomalies in the zonal-mean
541 zonal wind at the surface are most likely independent of the signals in the middle atmosphere. The alleged surface
542 solar signals in MSLP seem to mimic AO-positive (and AO-negative) patterns rather randomly than in a systematic
543 way. This applies to the HadSLP2 data (Figure 6) and to the model data (Figure 7), which both depict most pronounced
544 an AO-positive pattern in January and February at different lag years however in months, where the strong
545 stratospheric influence (in December) is already weak or even reverses sign in the model (compare Figure 4). With
546 respect to the suggested synchronization between the decadal solar forcing and the NAO (e.g., Thiéblemont et al.,
547 2015) we cannot find any meaningful relations in the MiKlip historical simulations. This is supported by the fact that
548 all ensemble members show very individual phase relations (i.e., positive/negative correlations and maximizations
549 during different lag years) between the solar and the NAO time series. Additionally, more robust correlations could
550 not be achieved in different time segments (i.e., periods with stronger or weaker solar forcing). These findings apply
551 to the seasonal winter mean (DJF) as well as to individual winter months (December, January and February). As a

552 consequence, the detected phase relations in the HadSLP2 dataset should be interpreted carefully with respect to
553 potential physical connections between the solar forcing and the NAO, in particular since the observations represent
554 only one single ensemble member.

555 In summary, we draw four major conclusions:

- 556 1. The decadal variations of the TST in the MiKlip historical simulations are a product of the 11-year solar
557 cycle. In the course of this, an increase in the solar intensity leads to enhanced radiative shortwave heating
558 rates and a warming of the TST. These findings are consistent with other modeling studies concerning the
559 imprints of the 11-year solar cycle in the tropical upper stratosphere (Matthes et al., 2004, 2006; Schmidt et
560 al., 2010; Ineson et al., 2011; Chiodo et al., 2012; Langematz et al., 2013). The solar signals in the TST are
561 statistically significant and robust and were detected by our correlation and MLR analyses.
- 562 2. The dynamical response of the NH during winter in the middle atmosphere shows a weak strengthening of
563 the polar vortex during solar maximum in the ensemble mean in the MLR analysis. However, the signals
564 (especially in the zonal-mean zonal wind) are mostly insignificant and of opposite sign in individual ensemble
565 members, and thus not a robust feature. We suppose that the dynamical background state in the middle
566 atmosphere (i.e., the variability of the polar vortex) seems to play an important role for the transfer of the
567 initial radiative solar signal from the upper tropical stratosphere down to the troposphere in NH winter. The
568 important role of middle atmosphere dynamics in modulating potential solar signals is currently investigated
569 as part of the SOLCHECK project and will be published in a subsequent paper (Wenjuan Huo, personal
570 communication).
- 571 3. The detected anomalies in the zonal-mean zonal wind and MSLP at the surface seem not to be related to the
572 timing of the seasonal march of the signals in the middle atmosphere and are most likely a manifestation of
573 the internal variability in the troposphere itself.
- 574 4. Concerning the decadal variations of the NAO and the solar forcing, our results suggest that both are
575 independent from each other. We find a range of phase relations between the NAO and the solar forcing
576 throughout our ensemble members, which implies a random statistical relation rather than a physical sound
577 connection.

578

579 It should be noted that we did not explicitly analyze a potential TSI controlled bottom-up effect on the solar surface
580 signal, as bottom-up effects are rather confined to tropical latitudes with a prolonged influence of the TSI throughout
581 the year (e.g., Meehl et al., 2008). Moreover, potential effects related to energetic particle precipitation are not
582 explicitly included in the MiKlip experiments. Since these effects are known to be rather small and even less
583 understood than the 11-year solar cycle surface imprints, we don't think they would alter our results significantly (see
584 the introduction section).

585 Since the critical study of Chiodo et al. (2019), the “top-down” mechanism and its surface imprints have been further
586 discussed in the scientific community. It is unquestionable that early studies with GCMs and CCMs found evidence
587 of a “top-down” mechanism in the middle atmosphere which in most cases penetrated into the troposphere in NH
588 winter (Matthes et al., 2004, 2006; Schmidt et al., 2010; Ineson et al., 2011; Chiodo et al., 2012; Langematz et al.,
589 2013). These studies all reproduced more or less the basic features of the “top-down” mechanism, thus confirming the
590 physical mechanisms at work suggested by Kodera and Kodera (2002). In contrast, more recent simulations with
591 CCMs and ESMs do not seem to find statistical responses of surface variables to the decadal solar forcing (e.g., Chiodo
592 et al., 2019; this study). Only Drews et al. (2022) showed a near-surface solar imprint for solar cycles with strong
593 amplitudes. The MiKlip simulations are more in line with Chiodo et al. (2019), who argued that the alleged surface
594 solar signals could be an incidental product which is only detectable during phases with stronger solar cycles. Our
595 results even suggest that robust solar surface imprints are basically absent throughout the complete historical period
596 and are thus not sensitive to the amplitude of individual solar cycles. At this point we would like to emphasize that in
597 contrast to previous studies, the MiKlip simulations represent a transient climate system driven by a realistic (observed)
598 solar forcing thus enhancing the confidence in a comparison of our model results to observations.

599 We suggest that the gradual ‘fading away’ of significant solar near-surface signatures in more up-to-date model studies
600 is closely related to progresses made in model development and computer capacities allowing for ensemble
601 simulations. The early simulations were conducted with fixed lower boundary conditions (i.e., prescribed SSTs from
602 observations or control run experiments) (Matthes et al., 2006; Schmidt et al., 2010; Chiodo et al., 2012). Some applied
603 perpetual conditions for the solar forcing (i.e., perpetual solar maximum vs. perpetual solar minimum) and steady-
604 state conditions for the greenhouse gas forcing (Matthes et al., 2006; Schmidt et al., 2010; Ineson et al., 2011). While
605 these models included the necessary physical mechanisms, i.e., UV radiation codes and middle atmosphere dynamics,
606 to capture the solar UV-induced top-down solar signal, the complex nature of physical and chemical processes and the

607 spectrum of internal variability were reduced. Prescribed SSTs, for example, prevent the model from developing the
608 complete spectrum of interannual variability in the troposphere (e.g., induced by the internal variability of the NAO),
609 which might counteract potential surface solar signals. In addition, steady-state background conditions in atmospheric
610 greenhouse gas concentrations and prescribed ozone depleting substances do not take into account transient adjustment
611 processes in the atmospheric dynamics, which again lead to a reduction of the overall internal variability and maybe
612 an overestimation of solar-induced signals. Moreover, due to more limited computer capacities, the results from the
613 early model studies were mostly based on single simulations.

614 In contrast, our results show that in a state-of-the-art climate model system the potential solar near-surface signals are
615 rather weak, not robust and inconsistent with the timing in the middle atmosphere. One potential reason is the
616 additional variability component introduced into the model by the interactively coupled ocean model. Misios and
617 Schmidt (2012) also showed the impact of an interactive ocean on the simulated solar response in the tropical Pacific
618 region. While individual ensemble simulations produce the expected phase correlation between the NAO and the solar
619 cycle, others show the opposite behavior. Thus, we do not find any convincing evidence in our model simulations of
620 the alleged decadal synchronization between the NAO and the solar forcing, as suggested by Thiéblemont et al. (2015).
621 In our view, the decadal near-surface signals detected in the MiKlip historical simulations are a product of the internal
622 variability in the troposphere itself and not a physical consequence of the “top-down” mechanism.

623 We would further like to mention that a strong reduction of the interannual variability in two basically independent
624 time series – be it by bandpass filtering like in our study or in Thiéblemont et al. (2015), or by using wide running
625 mean windows like in Drews et al. (2022) – will always lead to significant alignments of these two time series at some
626 point, if they are shifted towards each other gradually. Thus, the phase relations in our (and other studies) seem to be
627 a statistical artifact and not the consequence of a physical phase coupling. We also would like to question if the oceanic
628 memory is sensitive enough to store the tiny surface solar signals (even if there are some) for the duration of a complete
629 decade. Hence, in our opinion a much more profound solar forcing would be needed to significantly influence the
630 ocean temperature and thus dynamically driven feedbacks. Such forcings, however, typically operate on the centennial
631 timescale which is characterized by phases of Grands Solar Maxima and Minima (e.g., Spiegl and Langematz (2020)).
632 Also, please keep in mind the strong variability of the main pressure systems in the North Atlantic, which might wipe
633 out potential surface solar signals within a couple of months. The discrepancies between the observed and modelled
634 internal variability in response to external forcings (such as solar variability) may also be attributed to the "signal-to-

635 noise" paradox which states that relatively small changes in the external forcing will not lead to detectable changes in
636 the variability spectrum in both the real climate system and model simulations as discussed by Scaife and Smith (2018).
637 Furthermore, and in our opinion, a physically sound explanation for the alleged NAO-solar cycle phase coupling is
638 missing so far. Thus, the claim that an inclusion of the 11-year solar cycle would lead to a better understanding of the
639 decadal oscillations in the NH troposphere during winter, is not supported by our analyses of the MiKlip historical
640 ensemble simulations. Future studies with a distinct focus on the decadal prediction skill might help to confirm our
641 results.

642

643 **Data availability**

644 The main numerical results will be made available upon request by the authors.

645 **Author contributions**

646 TS was in charge in conducting the analysis and writing the manuscript. UL initiated the study and contributed to
647 writing the manuscript. HP and JK were involved in conducting the MiKlip historical simulations and writing the
648 manuscript

649 **Competing interests**

650 The authors do not declare any competing interests.

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660 **References**

- 661 Allan, R., and Ansell, T.: A new globally complete monthly historical gridded mean sea level pressure dataset
662 (HadSLP2): 1850–2004. *J. Climate*, 19(22), 5816-5842, doi.org/10.1175/JCLI3937.1, 2006.
- 663 Andrews, D. G.: Wave–mean-flow interaction in the middle atmosphere. In *Adv. Geophys.* (Vol. 28, pp. 249-275).
664 Elsevier, [doi.org/10.1016/S0065-2687\(08\)60226-5](https://doi.org/10.1016/S0065-2687(08)60226-5), 1985.
- 665 Andrews, M. B., Knight, J. R., and Gray, L. J.: A simulated lagged response of the North Atlantic Oscillation to the
666 solar cycle over the period 1960–2009. *Environ. Res. Lett.*, 10(5), 054022, doi.org/10.1088/1748-9326/10/5/054022,
667 2015.
- 668 Arsenovic, P., Rozanov, E., Stenke, A., Funke, B., Wissing, J. M., Mursula, K., ... and Peter, T.: The influence of
669 Middle Range Energy electrons on atmospheric chemistry and regional climate. *J Atmos Sol Terr Phys.*, 149,
670 180-190, doi.org/10.1016/j.jastp.2016.04.008, 2016.
- 671 Baldwin, M. P., and Dunkerton, T. J.: Stratospheric harbingers of anomalous weather regimes. *Science*, 294(5542),
672 581-584. doi.org/10.1126/science.1063315, 2001.
- 673 Baumgaertner, A. J., Jöckel, P., Riede, H., Stiller, G. and Funke, B.: Energetic particle precipitation in
674 ECHAM5/MESy–Part 2: Solar proton events. *Atmos. Chem. Phys.*, 10(15), 7285-7302, [doi.org/10.5194/acp-](https://doi.org/10.5194/acp-10-7285-2010)
675 [10-7285-2010](https://doi.org/10.5194/acp-10-7285-2010), 2010.
- 676 Bodeker, G. E., Boyd, I. S., and Matthews, W. A.: Trends and variability in vertical ozone and temperature profiles
677 measured by ozonesondes at Lauder, New Zealand: 1986–1996. *J. Geophys. Res.*, 103(D22), 28661-28681,
678 doi.org/10.1029/98JD02581, 1998.
- 679 Butchart, N.: The Brewer-Dobson circulation. *Rev. Geophys.*, 52(2), 157-184,
680 <https://doi.org/10.1002/2013RG000448>, 2014.
- 681 Cagnazzo, C., Manzini, E., Giorgetta, M. A., Forster, P. M. D. F., and Morcrette, J. J.: Impact of an improved
682 shortwave radiation scheme in the MAECHAM5 General Circulation Model, *Atmos. Chem. Phys.*, 7, 2503–
683 2515, <https://doi.org/10.5194/acp-7-2503-2007>, 2007.

684 Cionni, I., Eyring, V., Lamarque, J. F., Randel, W. J., Stevenson, D. S., Wu, F., ... and Waugh, D. W.: Ozone
685 database in support of CMIP5 simulations: results and corresponding radiative forcing. *Atmos. Chem. Phys.*,
686 11(21), 11267-11292, doi.org/10.5194/acp-11-11267-2011, 2011.

687 Chiodo, G., Calvo, N., Marsh, D. R., and Garcia-Herrera, R.: The 11 year solar cycle signal in transient simulations
688 from the Whole Atmosphere Community Climate Model. *J. Geophys. Res.*, 117(D6),
689 doi.org/10.1029/2011JD016393, 2012.

690 Chiodo, G., Oehrlein, J., Polvani, L. M., Fyfe, J. C., and Smith, A. K.: Insignificant influence of the 11-year solar
691 cycle on the North Atlantic Oscillation. *Nat. Geosci.*, 12(2), 94-99, doi.org/10.1038/s41561-018-0293-3, 2019.

692 Drews, A., Huo, W., Matthes, K., Kodera, K., and Kruschke, T.: The Sun's role in decadal climate predictability in
693 the North Atlantic. *Atmos. Chem. Phys.*, 22(12), 7893-7904, doi.org/10.5194/acp-22-7893-2022, 2022.

694 Eyring, V., Lamarque, J. F., Hess, P., Arfeuille, F., Bowman, K., Chipperfield, M. P., ... and Young, P. Y.: Overview
695 of IGAC/SPARC Chemistry-Climate Model Initiative (CCMI) community simulations in support of upcoming
696 ozone and climate assessments. *SPARC newsletter*, 40 (Januar), 48-66, 2013.

697 Forster, P. M., et al.: Evaluation of radiation scheme performance within chemistry climate models, *J. Geophys.*
698 *Res.*, 116, D10302, [doi:10.1029/2010JD015361](https://doi.org/10.1029/2010JD015361), 2011.

699 Fouquart, Y., and Bonnel, B.: Computations of solar heating of the earth's atmosphere—A new parameterization,
700 *Beitr. Phy. Atmos.*, 53, 35–62, 1980.

701 Gray, L. J., et al.: Solar influences on climate, *Rev. Geophys.*, 48, RG4001, [doi:10.1029/2009RG000282](https://doi.org/10.1029/2009RG000282), 2010.

702 Gray, L. J., Scaife, A. A., Mitchell, D. M., Osprey, S., Ineson, S., Hardiman, S., ... and Kodera, K.: A lagged
703 response to the 11 year solar cycle in observed winter Atlantic/European weather patterns. *J. Geophys. Res.*,
704 118(24), 13-405, doi.org/10.1002/2013JD020062, 2013.

705 Huang, J., Hitchcock, P., Maycock, A. C., McKenna, C. M., and Tian, W.: Northern hemisphere cold air outbreaks
706 are more likely to be severe during weak polar vortex conditions. *Communications Earth and Environment*,
707 2(1), 147. doi.org/10.1038/s43247-021-00215-6, 2021.

708 Iacono, M. J., Delamere, J.S., Mlawer, E.J., Shephard, M. W., Clough, S. A., and Collins, W. D.: Radiative forcing
709 by long-lived greenhouse gases: Calculations with the AER radiative transfer models, *J. Geophys. Res.*, 113,
710 D13103, [doi:10.1029/2008JD009944](https://doi.org/10.1029/2008JD009944), 2008.

711 Ilyina, T., Six, K. D., Segsneider, J., Maier-Reimer, E., Li, H., and Núñez-Riboni, I.: Global ocean
712 biogeochemistry model HAMOCC: Model architecture and performance as component of the MPI-Earth system
713 model in different CMIP5 experimental realizations. *J. Adv. Model. Earth Syst.*, 5(2), 287-315,
714 doi.org/10.1029/2012MS000178, 2013.

715 Ineson, S., Scaife, A. A., Knight, J. R., Manners, J. C., Dunstone, N. J., Gray, L. J., and Haigh, J. D.: Solar forcing of
716 winter climate variability in the Northern Hemisphere. *Nat. Geosci.*, 4(11), 753-757, doi.org/10.1038/ngeo1282,
717 2011.

718 Jackman, C. H., Marsh, D. R., Vitt, F. M., Garcia, R. R., Randall, C. E., Fleming, E. L. and Frith, S. M.: Long-term
719 middle atmospheric influence of very large solar proton events. *J. Geophys. Res.*, 114(D11),
720 doi.org/10.1029/2008JD011415, 2009.

721 Jungclaus, J. H., Fischer, N., Haak, H., Lohmann, K., Marotzke, J., Matei, D., ... and Von Storch, J. S.:
722 Characteristics of the ocean simulations in the Max Planck Institute Ocean Model (MPIOM) the ocean
723 component of the MPI-Earth system model. *J. Adv. Model. Earth Syst.*, 5(2), 422-446,
724 doi.org/10.1002/jame.20023, 2013.

725 Kodera, K.: Solar cycle modulation of the North Atlantic Oscillation: Implication in the spatial structure of the
726 NAO. *Geophys. Res. Lett.*, 29(8), 59-1, doi.org/10.1029/2001GL014557, 2002.

727 Kodera, K., and Kuroda, Y.: Dynamical response to the solar cycle. *J. Geophys. Res.*, 107(D24),
728 doi.org/10.1029/2002JD002224, ACL-5. 2002.

729 Kuroda, Y., Kodera, K., Yoshida, K., Yukimoto, S., & Gray, L.: Influence of the solar cycle on the North Atlantic
730 Oscillation. *J. Geophys. Res.*, 127(1), e2021JD035519, doi.org/10.1029/2021JD035519, 2022.

731 Langematz, U., Kubin A., Brühl, C., Baumgaertner, A.J.G., Cubasch, U., and Spangehl, T.: Solar effects on
732 chemistry and climate including ocean interactions, in *Climate And Weather of the Sun-Earth System*
733 (CAWSES): Highlights from a Priority Program, F.-J. Lübken, ed., Springer, Dordrecht, The Netherlands, 2013.

734 Lean, J.: Evolution of the Sun's spectral irradiance since the Maunder Minimum. *Geophys. Res. Lett.*, 27(16), 2425-
735 2428, doi.org/10.1029/2000GL000043, 2000.

736 Ma, H., Chen, H., Gray, L., Zhou, L., Li, X., Wang, R. and Zhu, S.: Changing response of the North
737 Atlantic/European winter climate to the 11-year solar cycle. *Environ. Res. Lett.*, 13(3), 034007,
738 doi.org/10.1088/1748-9326/aa9e94, 2018.

739 Marotzke, J., Müller, W. A., Vamborg, F. S., Becker, P., Cubasch, U., Feldmann, H., ... and Ziese, M.: MiKlip: a
740 national research project on decadal climate prediction. *Bull. Am. Meteorol. Soc.*, 97(12), 2379-2394,
741 <https://doi.org/10.1175/BAMS-D-15-00184.1>, 2016.

742 Matthes, K., Langematz, U., Gray, L. L., Kodera, K., and Labitzke, K.: Improved 11-year solar signal in the Freie
743 Universität Berlin climate middle atmosphere model (FUB-CMAM). *J. Geophys. Res.*, 109(D6),
744 doi.org/10.1029/2003JD004012, 2004.

745 Matthes, K., Kuroda, Y., Kodera, K., and Langematz, U.: Transfer of the solar signal from the stratosphere to the
746 troposphere: Northern winter. *J. Geophys. Res.*, 111(D6), doi.org/10.1029/2005JD006283, 2006.

747 Matthes, K., Funke, B., Andersson, M. E., Barnard, L., Beer, J., Charbonneau, P., Clilverd, M. A., Dudok de Wit, T.,
748 Haberreiter, M., Hendry, A., Jackman, C. H., Kretzschmar, M., Kruschke, T., Kunze, M., Langematz, U.,
749 Marsh, D. R., Maycock, A. C., Misios, S., Rodger, C. J., Scaife, A. A., Seppälä, A., Shangguan, M., Sinnhuber,
750 M., Tourpali, K., Usoskin, I., van de Kamp, M., Verronen, P. T., and Versick, S.: Solar forcing for CMIP6
751 (v3.2), *Geosci. Model Dev.*, 10, 2247–2302, doi.org/10.5194/gmd-10-2247-2017, 2017.

752 Meehl, G. A., J. M. Arblaster, G. Branstator, and H. van Loon: A Coupled Air–Sea Response Mechanism to Solar
753 Forcing in the Pacific Region. *J. Climate*, 21, 2883–2897, <https://doi.org/10.1175/2007JCLI1776.1>, 2008.

754 Meehl, G. A., Goddard, L., Boer, G., Burgman, R., Branstator, G., Cassou, C., ... and Yeager, S.: Decadal climate
755 prediction: an update from the trenches. *Bull. Am. Meteorol. Soc.*, 95(2), 243-267, doi.org/10.1175/BAMS-D-12-00241.1, 2014.

757 Mehta, V., Meehl, G., Goddard, L., Knight, J., Kumar, A., Latif, M., ... and Stammer, D.: Decadal climate
758 predictability and prediction: where are we?. *Bull. Am. Meteorol. Soc.*, 92(5), 637-640,
759 doi.org/10.1175/2010BAMS3025.1, 2011.

760 Misios, S. and Schmidt, H.: Mechanisms Involved in the Amplification of the 11-yr solar cycle signal in the tropical
761 Pacific Ocean, *J. Climate*, 25, 5102–5118, doi.org/10.1175/JCLID-11-00261.1, 2012.

762 Müller, W. A., Jungclaus, J. H., Mauritsen, T., Baehr, J., Bittner, M., Budich, R., ... and Marotzke, J: A higher-
763 resolution version of the max planck institute earth system model (MPI-ESM1. 2-HR). *J. Adv. Model. Earth*
764 *Syst.*, 10(7), 1383-1413, doi.org/10.1029/2017MS001217, 2018.

765 Paulsen, H., Ilyina, T., Six, K. D., and Stemmler, I.: Incorporating a prognostic representation of marine nitrogen
766 fixers into the global ocean biogeochemical model HAMOCC. *J. Adv. Model. Earth Syst.*, 9(1), 438-464,
767 doi.org/10.1002/2016MS000737, 2017.

768 Pohlmann, H., Müller, W. A., Kulkarni, K., Kameswarrao, M., Matei, D., Vamborg, F. S. E., Kadow, C., Illing, S.,
769 and Marotzke, J.: Improved forecast skill in the tropics in the new MiKlip decadal climate predictions, *Geophys.*
770 *Res. Lett.*, 40, 5798– 5802, [doi.10.1002/2013GL058051](https://doi.org/10.1002/2013GL058051), 2013.

771 Pohlmann, H., Müller, W. A., Bittner, M., Hettrich, S., Modali, K., Pankatz, K., and Marotzke, J.: Realistic quasi-
772 biennial oscillation variability in historical and decadal hindcast simulations using CMIP6 forcing. *Geophys.*
773 *Res. Lett.*, 46(23), 14118-14125, doi.org/10.1029/2019GL084878, 2019.

774 Randel, W. J., Smith, A. K., Wu, F., Zou, C., and Qian, H.: Stratospheric Temperature Trends over 1979–2015
775 Derived from Combined SSU, MLS, and SABER Satellite Observations, *J. Climate*, 29(13), 4843-4859,
776 doi.org/10.1175/JCLI-D-15-0629.1, 2016.

777 Reick, C. H., Raddatz, T., Brovkin, V., and Gayler, V.: Representation of natural and anthropogenic land cover
778 change in MPI-ESM. *J. Adv. Model. Earth Syst.*, 5(3), 459-482, doi.org/10.1002/jame.20022, 2013.

779 Scaife, A. A., Ineson, S., Knight, J. R., Gray, L., Kodera, K., and Smith, D. M.: A mechanism for lagged North
780 Atlantic climate response to solar variability. *Geophys. Res. Lett.*, 40(2), 434-439, doi.org/10.1002/grl.50099,
781 2013.

782 Scaife, A. A. and Smith, D.: A signal-to-noise paradox in climate science. *NPJ Climate and Atmospheric Science*,
783 1(1), 28, [doi:10.1038/s41612-018-0038-4](https://doi.org/10.1038/s41612-018-0038-4), 2018.

784 Seppälä, A., Randall, C. E., Clilverd, M. A., Rozanov, E., & Rodger, C. J.: Geomagnetic activity and polar surface
785 air temperature variability. *J. Geophys. Res.*, 114(A10), doi.org/10.1029/2008JA014029, 2009.

786 Seppälä, A., Matthes, K., Randall, C. E. and Mironova, I. A.: What is the solar influence on climate? Overview of
787 activities during CAWSES-II. *Progress in Earth and Planetary Science*, 1(1), 1-12, [doi.org/10.1186/s40645-014-](https://doi.org/10.1186/s40645-014-0024-3)
788 [0024-3](https://doi.org/10.1186/s40645-014-0024-3), 2014.

789 Schmidt, H., Brasseur, G. P., and Giorgetta, M. A.: Solar cycle signal in a general circulation and chemistry model
790 with internally generated quasi-biennial oscillation. *J. Geophys. Res.*,115(D1), doi.org/10.1029/2009JD012542,
791 2010.

792 Spiegel, T. and Langematz, U.: Twenty-First-Century Climate Change Hot Spots in the Light of a Weakening Sun.
793 *Journal of Climate*, 33(9), 3431-3447, doi.org/10.1175/JCLI-D-19-0059.1, 2020.

794 Stevens, B., et al.: Atmospheric component of the MPI-M Earth System Model: ECHAM6, *J. Adv. Model. Earth*
795 *Syst.*, 5, 146–172, [doi:10.1002/jame.20015](https://doi.org/10.1002/jame.20015), 2013.

796 Taylor, K. E., Stouffer, R. J., and Meehl, G. A.: An overview of CMIP5 and the experiment design. *Bull. Am.*
797 *Meteorol. Soc.*, 93(4), 485-498. doi.org/10.1175/BAMS-D-11-00094.1, 2012.

798 Thiéblemont, R., Matthes, K., Omrani, N. E., Kodera, K., and Hansen, F.: Solar forcing synchronizes decadal North
799 Atlantic climate variability. *Nat. Commun.*, 6(1), 8268, [doi:10.1038/ncomms9268](https://doi.org/10.1038/ncomms9268), 2015.

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